

*Citation:*

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**Mathematics.** — “Homogeneous linear differential equations of order two with given relation between two particular integrals.”  
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 (5<sup>th</sup> communication).

(Communicated in the meeting of April 26, 1912).

The equations (8) and (29) (see 1<sup>st</sup> comm. p. 393 and 398) show us in the case that the equation  $F(x, y, z) = 0$  represents a conic (see for the notation: 4<sup>th</sup> comm. p. 1015):

$$q_1 = \frac{c^2 z^3 H}{(n-1)^2 F_z^2} = \frac{\Delta z^3}{g^2} = e^{\int I dt},$$

where  $c$  is put equal to 1.

From this ensues

$$\frac{\dot{g}}{g} = -\frac{I}{3}.$$

Let us further put:

$$g = a_{11} \zeta^2 z, \dots \dots \dots (72)$$

we then find:

$$\frac{\dot{\zeta}}{\zeta} = \frac{1}{2} \frac{\dot{g}}{g} = -\frac{I}{6},$$

or

$$I = -6 \frac{\dot{\zeta}}{\zeta} \dots \dots \dots (73)$$

The equation (62) (see 4<sup>th</sup> comm. p. 1015) runs now as follows:

$$I^2 = 36 \frac{\dot{\zeta}^2}{\zeta^2} = \frac{9}{a_{11} \Delta z^2 \zeta^2} (-a_{11}^2 A_{11} z^2 \zeta^4 + 2a_{11} \Delta z^2 \zeta^2 - a_{11} \Delta z^2),$$

or making use of the notation (59) (4<sup>th</sup> comm. p. 1003),

$$4\zeta^2 = -\lambda^2 \zeta^4 + 2\zeta^2 - 1; \dots \dots \dots (74)$$

so  $\zeta$  is likewise an elliptic function of  $\tau$ . Its invariant has the same value (68) as that of the function  $u = I^2$  (compare (67) <sup>1)</sup> (4<sup>th</sup> comm. p. 1006).

We can now deduce out of the equation

$$A_{11}x - A_{11}y = \sqrt{-A_{11}g^2 + 2\Delta gz - \Delta a_{11}z^2} = \sqrt{a_{11}\Delta} \cdot z \sqrt{-\lambda^2 \zeta^4 + 2\zeta^2 - 1} \quad (75)$$

(see 4<sup>th</sup> comm. p. 1005 at the bottom)

$$A_{11}x - A_{11}y = 2z\sqrt{a_{11}\Delta} \cdot \zeta \dots \dots \dots (76)$$

<sup>1)</sup> In the 4<sup>th</sup> comm. in the table on p. 1014 and in the enumeration of the cases on p. 1015  $\delta_1 = e^{i\psi}$  and  $\delta_2 = e^{-i\psi}$  must be replaced by  $\delta_1 = e^{-i\psi}$ ,  $\delta_2 = e^{+i\psi}$ .

As from (73) follows

$$a_{11}x + a_{21}y = a_{31}z(\zeta^2 - 1), \dots \dots \dots (77)$$

we find with the aid of (76) and (77)

$$\left. \begin{aligned} (a_{11}A_{11} + a_{21}A_{21})x &= (\Delta - a_{31}A_{31})x = \Delta(1 - \lambda^2)x = \\ &= \{2a_{21}\sqrt{a_{31}\Delta} \cdot \zeta + a_{31}A_{31}(\zeta^2 - 1)\}z, \\ (a_{11}A_{11} + a_{21}A_{21})y &= (\Delta - a_{31}A_{31})y = \Delta(1 - \lambda^2)y = \\ &= \{-2a_{11}\sqrt{a_{31}\Delta} \cdot \zeta + a_{31}A_{31}(\zeta^2 - 1)\}z. \end{aligned} \right\} (78)$$

In this way we have expressed  $x$  and  $y$  as functions of  $\tau$  with the aid of the function  $\zeta$ . It is now still our task to determine  $\zeta$  as function of  $\tau$ . Let us now put in

$$9\dot{u}^2 = u^2 - 36u^2 + 324(1 - \lambda^2)u$$

(see 4<sup>th</sup> comm. p. 1016)

$$u = I^2 = 36v + 12, \dots \dots \dots (79)$$

we then find

$$\dot{v}^2 = 4v^3 - \frac{1 + 3\lambda^2}{3}v - \frac{9\lambda^2 - 1}{27}.$$

By applying the ordinary notation

$$\frac{1 + 3\lambda^2}{3} = g_2, \quad \frac{9\lambda^2 - 1}{27} = g_3, \dots \dots \dots (80)$$

we then find

$$v = p(\tau; g_2, g_3)$$

and

$$I = \pm 6 \sqrt{p(\tau; g_2, g_3) + \frac{1}{3}}, \dots \dots \dots (81)$$

so that

$$\frac{\zeta}{\xi} = \mp \sqrt{p(\tau; g_2, g_3) + \frac{1}{3}} \dots \dots \dots (82)$$

Before transforming the  $p$ -function of WEIERSTRASS we wish to remind the readers that the roots of  $\dot{u} = 0$  are

$$u_1 = 0, \quad u_2 = 18(1 + \lambda), \quad u_3 = 18(1 - \lambda),$$

so that for the roots of  $\dot{v} = 0$  (see (79)) we find

$$v_1 = -\frac{1}{3}, \quad v_2 = \frac{1 + 3\lambda}{6}, \quad v_3 = \frac{1 - 3\lambda}{6}.$$

We shall now investigate the relative value of these roots in the three cases: II ( $+1 < \lambda < +\infty$ ), IV ( $+1 > \lambda > 0$ ), VI ( $\lambda = i\lambda'$ ) (see 4<sup>th</sup> comm. p. 1014).

Case II:  $+1 < \lambda < +\infty$

$$v_1 = -\frac{1}{3}, v_2 > \frac{2}{3}, v_3 < -\frac{1}{3}.$$

The roots are all real. Let us call them in the ordinary way in descending order  $e_1, e_2, e_3$ , we then find

$$e_1 = \frac{1+3\lambda}{6}, e_2 = -\frac{1}{3}, e_3 = \frac{1-3\lambda}{6} \dots \dots \dots II$$

Case IV:  $+1 > \lambda > 0$ .

$$v_1 = -\frac{1}{3}, \frac{2}{3} > v_2 > \frac{1}{6}, \frac{1}{6} > v_3 > -\frac{1}{3}.$$

The roots are here, too, all real and run when arranged:

$$e_1 = \frac{1+3\lambda}{6}, e_2 = \frac{1-3\lambda}{6}, e_3 = -\frac{1}{3} \dots \dots \dots IV$$

Case VI:  $\lambda = i\lambda'$ .

The roots  $v_2$  and  $v_3$  are now conjugate complex. If we follow the notation generally assumed, we then write:

$$e_2' = -\frac{1}{3}, e_1' = \frac{1+3i\lambda'}{6}, e_3' = \frac{1-3i\lambda'}{6} \dots \dots \dots VI$$

When reducing the  $p$ -functions to the elliptic functions of JACOBI we make use of the following formulae of reduction: <sup>1)</sup>

$$sn(v) = \sqrt{\frac{e_1 - e_3}{p(\tau) - e_3}}, \quad cn(v) = \sqrt{\frac{p(\tau) - e_1}{p(\tau) - e_3}}, \quad dn(v) = \sqrt{\frac{p(\tau) - e_2}{p(\tau) - e_3}},$$

$$v = \tau \sqrt{e_1 - e_3}, \quad k^2 = \frac{e_2 - e_3}{e_1 - e_3}, \quad k'^2 = \frac{e_1 - e_2}{e_1 - e_3};$$

$$p(\tau; e_1', e_2', e_3') = e_2' + \frac{e_1' - e_3'}{4ikk'} \cdot \frac{cn^2(v)}{sn^2(v) \cdot dn^2(v)},$$

$$v = \tau \sqrt{(e_2' - e_3')(e_2' - e_1')}, \quad k^2 = \frac{-3e_2' + 2\sqrt{(e_2' - e_3')(e_2' - e_1')}}{4\sqrt{(e_2' - e_3')(e_2' - e_1')}},$$

$$k'^2 = \frac{+3e_2' + 2\sqrt{(e_2' - e_3')(e_2' - e_1')}}{4\sqrt{(e_2' - e_3')(e_2' - e_1')}}}, \quad kk' = \frac{\sqrt{-9e_2'^2 + 4(e_2' - e_3')(e_2' - e_1')}}{4\sqrt{(e_2' - e_3')(e_2' - e_1')}}}.$$

The expression for  $\xi : \zeta$  becomes in this way:

in case II

$$\frac{\xi}{\zeta} = \mp \sqrt{p(\tau; g_2, g_3) - e_3} = \mp \sqrt{e_1 - e_3} \cdot \frac{dn(v)}{sn(v)} \left. \vphantom{\frac{\xi}{\zeta}} \right\} v = \tau \sqrt{e_1 - e_3},$$

in case IV

$$\frac{\xi}{\zeta} = \mp \sqrt{p(\tau; g_2, g_3) - e_3} = \pm \sqrt{e_1 - e_3} \cdot \frac{1}{sn(v)} \left. \vphantom{\frac{\xi}{\zeta}} \right\} k^2 = \frac{e_2 - e_3}{e_1 - e_3}; \quad k'^2 = \frac{e_1 - e_2}{e_1 - e_3}$$

<sup>1)</sup> See i. a. M. KRAUSE: Theorie der elliptischen Funktionen (Leipzig, TEUBNER (p. 135, 136, 147, 148).

in case VI

$$\frac{\dot{\zeta}}{\zeta} = \mp \sqrt{p(\tau; e_1', e_2', e_3') - e_2'} = \mp \sqrt{\frac{e_1' - e_3'}{4 i k k'} \cdot \frac{cn(v)}{sn(v) \cdot dn(v)}}$$

$$v = \tau \sqrt{(e_2' - e_3')(e_2' - e_1')}, \quad k k' = \frac{\sqrt{-9e_2'^2 + 4(e_2' - e_3')(e_2' - e_1')}}{\sqrt{(e_2' - e_3')(e_2' - e_1')}},$$

or, after having expressed the roots  $e_1, e_2, e_3, e_1', e_2', e_3'$  in  $\lambda$ :

in case II

$$\frac{\dot{\zeta}}{\zeta} = \mp \sqrt{\lambda} \cdot \frac{dn(v)}{sn(v)}; \quad v = \tau \sqrt{\lambda}, \quad k^2 = \frac{\lambda - 1}{2\lambda}, \quad k'^2 = \frac{\lambda + 1}{2\lambda};$$

in case IV

$$\frac{\dot{\zeta}}{\zeta} = \mp \sqrt{\frac{1 + \lambda}{2}} \cdot \frac{1}{sn(v)}; \quad v = \tau \sqrt{\frac{1 + \lambda}{2}}, \quad k^2 = \frac{1 - \lambda}{1 + \lambda}, \quad k'^2 = \frac{2\lambda}{1 + \lambda}; \quad (83)$$

in case VI

$$\frac{\dot{\zeta}}{\zeta} = \mp \sqrt{\frac{\sqrt{1 + \lambda'^2}}{2}} \cdot \frac{cn(v)}{sn(v) \cdot dn(v)}; \quad v = \tau \sqrt{\frac{\sqrt{1 + \lambda'^2}}{2}};$$

$$k^2 = \frac{1 + \sqrt{1 + \lambda'^2}}{2\sqrt{1 + \lambda'^2}}, \quad k'^2 = \frac{-1 + \sqrt{1 + \lambda'^2}}{2\sqrt{1 + \lambda'^2}}.$$

Let us substitute these expressions in (14), we then find successively

$$\text{in case II } \zeta_1 = \pm \frac{i}{\sqrt{\lambda}} \cdot \frac{1 + cn(v)}{sn(v)}, \quad \zeta_2 = \pm \frac{1}{\lambda \zeta_1},$$

$$\text{in case IV } \zeta_1 = \pm \frac{i}{\lambda} \sqrt{\frac{1 + \lambda}{2}} \cdot \frac{cn(v) + dn(v)}{sn(v)}, \quad \zeta_2 = \pm \frac{1}{\lambda \zeta_1},$$

$$\text{in case VI } \zeta_1 = \pm \frac{\sqrt{4(1 + \lambda'^2)}}{\lambda'} \cdot \frac{dn(v)}{sn(v)}, \quad \zeta_2 = \pm \frac{i}{\lambda' \zeta_1}.$$

Let us now choose

$$l = + 6 \sqrt{p(\tau; g_2, g_3) + \frac{1}{3}}$$

and for  $\zeta$  the expressions  $\zeta_1$  with the upper sign, we find:

$$\begin{aligned}
II \quad I &= +6\sqrt{\lambda} \cdot \frac{dn(v)}{sn(v)} ; v = \tau\sqrt{\lambda} , k^2 = \frac{\lambda-1}{2\lambda} , k'^2 = \frac{\lambda+1}{2\lambda} ; \\
\zeta &= \frac{+i}{\sqrt{\lambda}} \cdot \frac{1+cn(v)}{sn(v)} , \dot{\zeta} = -i \cdot \frac{\{1+cn(v)\}dn(v)}{sn^2(v)} . \\
IV \quad I &= +6 \sqrt{\frac{1+\lambda}{2}} \cdot \frac{1}{sn(v)} ; v = \tau \sqrt{\frac{1+\lambda}{2}} , \\
k^2 &= \frac{1-\lambda}{1+\lambda} , k'^2 = \frac{2\lambda}{1+\lambda} ; \\
\zeta &= + \frac{i}{\lambda} \sqrt{\frac{1+\lambda}{2}} \cdot \frac{cn(v)+dn(v)}{sn(v)} , \dot{\zeta} = \frac{-i(1+\lambda)}{2\lambda} \cdot \frac{cn(v)+dn(v)}{sn^2(v)} . \\
VI \quad I &= +6 \sqrt{\frac{\sqrt{1+\lambda'^2}}{2}} \cdot \frac{cn(v)}{sn(v) \cdot dn(v)} , v = \sqrt{\frac{\sqrt{1+\lambda'^2}}{2}} , \\
k^2 &= \frac{1+\sqrt{1+\lambda'^2}}{2\sqrt{1+\lambda'^2}} , k'^2 = \frac{-1+\sqrt{1+\lambda'^2}}{2\sqrt{1+\lambda'^2}} \\
\zeta &= \frac{\sqrt{4(1+\lambda'^2)}}{\lambda'} \cdot \frac{dn(v)}{sn(v)} , \dot{\zeta} = - \frac{\sqrt{1+\lambda'^2}}{\lambda'} \cdot \frac{cn(v)}{sn^2(v)} .
\end{aligned} \tag{84}$$

Let us restrict ourselves to real points  $(x, y)$  of the conic, then follows from (78) that  $\sqrt{a_{22}\Delta} \cdot \dot{\zeta}$  must always be real.

Case II (in which  $\lambda$  is real) appears only with the hyperbola for which holds  $A_{22} < 0$ ; so we have here

$$a_{22}\Delta = \frac{a_{22}A_{22}}{\Delta} \cdot \frac{\Delta^2}{A_{22}} = \frac{\lambda^2\Delta^2}{A_{22}} < 0.$$

From this ensues that in case II we shall find  $\dot{\zeta}$  always imaginary, and therefore  $\frac{\{1+cn(v)\}dn(v)}{sn^2(v)}$  is real;

Case IV is found with the hyperbola as well as with the ellipse. As here too  $\lambda$  is real we find

IVa. with the hyperbola ( $A_{22} < 0$ )  $a_{22}\Delta < 0$ , so  $\dot{\zeta}$  is imaginary or  $\frac{cn(v)+dn(v)}{sn^2(v)}$  real;

IVb. with the ellipse ( $A_{22} > 0$ )  $a_{22}\Delta > 0$ , so  $\zeta$  is real and  $\frac{cn(v)+dn(v)}{sn^2(v)}$  is purely imaginary.

Also case VI appears with the hyperbola as well as with the ellipse. On account of  $\lambda$  being purely imaginary, thus  $\lambda^2$  negative, holds:

VIa. for the hyperbola ( $A_{22} < 0$ )  $a_{22}\Delta > 0$ , hence  $\dot{\zeta}$  real, and  $\frac{cn(v)}{sn^2(v)}$  real;

VIb. for the ellipse ( $A_{22} > 0$ )  $a_{22}\Delta < 0$ , thus  $\zeta$  purely imaginary and also  $\frac{cn(v)}{sn^2(v)}$  purely imaginary.

From the preceding we see that  $v$  must move in its complex plane on the sides of the rectangles of the net formed by the lines  $v = mK +$  purely imaginary and  $v = niK' +$  real.

The value of  $\zeta^2 = \frac{g}{a_{22}} = \frac{a_{12}x + a_{22}y}{a_{22}} + 1$  is evidently positive on that side of the polar line  $g = 0$  of  $O$  with respect to the conic where  $O$  lies itself; on the other side  $\zeta^2$  is negative. The polar line  $g = 0$  of  $O$  divides therefore the plane into two parts: in one (in which  $O$  lies)  $\zeta$  is real, in the other  $\zeta$  is imaginary.

In the points of contact  $R_1$  and  $R_2$  of the tangents out of  $O$  to the conic  $\zeta$  is 0, so  $I = \infty$ .

In the points at infinity  $S_1$  and  $S_2$  we find that  $\zeta$  and  $\zeta$  are both infinite and  $I$  is also equal to  $\infty$ .

The diameter passing through  $O$  ( $A_{22}x - A_{12}y = 0$ ) intersects the conic in two points  $T_1$  and  $T_2$ , for which  $\zeta = 0$ , thus  $I = 0$ .

If we substitute the expressions (84) for  $\zeta$  and  $\zeta$  in the formulae (78) we at last arrive at  $x$  and  $y$  as functions of  $\tau$ .

With a view to  $\sqrt{A_{22}}$  being real or not, we shall deal with the cases of IV and VI separately. Furthermore we shall express  $\lambda$  everywhere

in  $\sigma = \frac{1-\lambda}{1+\lambda}$ , thus [in the anharmonic ratio of the four points

$R_1, R_2, S_1, S_2$ . We shall give the formulae for  $x$  only. The expressions for  $y$  we can easily find by replacing  $a_{22}$  in those for  $x$  by  $-a_{12}$  and  $A_{22}$  by  $A_{12}$ .

We then find at last:

$$II \quad x = \frac{1-\sigma}{2\sigma} \cdot \frac{1+cn(v)}{sn^2(v)} \left[ - (1+\sigma) \frac{a_{12}}{\sqrt{-A_{22}}} dn(v) + \frac{A_{12}}{-A_{22}} \{1+\sigma cn(v)\} \right],$$

$$v = \tau \sqrt{\frac{1-\sigma}{1+\sigma}};$$

$$IVa \quad x = \frac{1}{2\sigma} \cdot \frac{cn(v)+dn(v)}{sn^2(v)} \left[ - (1+\sigma) \frac{a_{22}}{\sqrt{-A_{22}}} + \frac{A_{12}}{-A_{22}} \{\sigma cn(v) + dn(v)\} \right],$$

$$v = \tau \sqrt{\frac{1}{1+\sigma}};$$

$$IVb \quad x = \frac{1}{2\sigma} \cdot \frac{cn(v)+dn(v)}{i sn^2(v)} \left[ - (1+\sigma) \frac{a_{22}}{\sqrt{-A_{22}}} - \frac{A_{12}}{A_{22}} i \{\sigma cn(v) + dn(v)\} \right],$$

$$v = \tau \sqrt{\frac{1}{1+\delta}};$$

$$VIIa \quad x = \frac{2 \operatorname{cn}(v)}{\operatorname{sn}^2(v)} \left[ \cos \frac{\psi}{2} \cdot \frac{a_{21}}{\sqrt{-A_{11}}} + 2 \cos^2 \frac{\psi}{4} \cdot \frac{A_{11}}{-A_{11}} \operatorname{cn}(v) \right],$$

$$v = \frac{\tau}{\sqrt{2 \cos \frac{\psi}{2}}}, \quad \delta = e^{-i\psi}, \quad \psi = i \log \delta;$$

$$VIIb \quad x = \frac{2 \operatorname{cn}(v)}{i \operatorname{sn}^2(v)} \left[ \cos \frac{\psi}{2} \cdot \frac{a_{21}}{\sqrt{A_{11}}} - 2 \cos^2 \frac{\psi}{4} \cdot \frac{A_{11}}{A_{11}} i \operatorname{cn}(v) \right],$$

$$v = \frac{\tau}{\sqrt{2 \cos \frac{\psi}{2}}}, \quad \delta = e^{-i\psi}, \quad \psi = i \log \delta.$$

When point  $(x, y)$  describes the conic, the variable  $v$  will describe a certain curve in its complex plane. This curve we shall investigate in the five cases mentioned above whilst at the same time we shall indicate how the functions  $\zeta$ ,  $\dot{\zeta}$  and  $I$  bear themselves during that motion.

*Case II.* Point  $O$  lies in the domain of the conjugate hyperbola; the diameter through  $O$  does not intersect the curve, i.e. the points  $T_1$  and  $T_2$  are imaginary. On the contrary the points  $R_1, R_2, S_1, S_2$  are all real.

	in $S_1^\infty$	on $S_1^\infty R_1$	in $R_1$	on $R_1 S_2^\infty$	in $S_2^\infty$	on $S_2^\infty R_2$	in $R_2$	on $R_2 S_1^\infty$	in $S_1^\infty$
$v$	0	purely imag.	$2iK'$	$2iK' + \text{real}$	$2K + 2iK'$	$2K + p.\text{imag}$	$2K$	real	0
$\zeta$	$\infty$	pos. real	0	pos. imag.	$\infty$	pos. real	0	pos. imag.	$\infty$
$\dot{\zeta}$		pos. imag.	$+\frac{i}{2}$	pos. imag.	$\infty$	neg. imag.	$-\frac{i}{2}$	neg. imag.	$\infty$
$I$	$\infty$	neg. imag.	$\infty$	neg. real	$\infty$	pos. imag.	$\infty$	pos. real	$\infty$

Here the curves are sketched which are described by  $v$  and  $I$  in their respective complex planes.

The points where  $I$  turns its direction of motion are arrived at by putting  $\dot{I} = 0$ . We then find the values of  $I$  corresponding to the roots of  $\dot{u} = 0$ ; these are  $u_1 = 0$ ,  $u_2 = \infty$ ,  $u_3 = 18(1+\lambda)$ ,  $u_4 = 18(1-\lambda)$ ; or  $I_1 = 0$ ,  $I_2 = \infty$ ,  $I_3 = 6\sqrt{\frac{1+\lambda}{2}}$ ,



IVa	in $S_1^\infty$	on $S_1^\infty R_1$	in $R_1$	on $R_1 T_1$	in $T_1$	on $T_1 R_2$	in $R_2$	on $R_2 S_2^\infty$	in $S_2^\infty$	on $S_2^\infty T_2$	in $T_2$	on $T_2 S_1^\infty$	in $S_1^\infty$
$\nu$	0	real	$2K$	$2K+p.\text{imag.}$	$2K+iK'$	$2K+p.\text{imag.}$	$2K+2iK'$	$2iK'+\text{real}$	$4K+2iK'$	$4K+p.\text{imag.}$	$4K+iK'$	$4K+p.\text{imag.}$	$4K$
$\zeta$	$\infty$	pos. imag.	0	pos. real	$\frac{\sqrt{1+\lambda}-\sqrt{1-\lambda}}{\lambda\sqrt{2}}$	pos. real	0	pos. imag.	$\infty$	pos. real	$\frac{\sqrt{1+\lambda}+\sqrt{1-\lambda}}{\lambda\sqrt{2}}$	pos. real	$\infty$
$\dot{\zeta}$	$\infty$	neg. imag.	$-\frac{i}{2}$	neg. imag.	0	pos. imag.	$+\frac{i}{2}$	pos. imag.	$\infty$	neg. imag.	0	pos. imag.	$\infty$
$I$	$\infty$	pos. real	$\infty$	pos. imag.	0	neg. imag.	$\infty$	neg. real	$\infty$	pos. imag.	0	neg. imag.	$\infty$

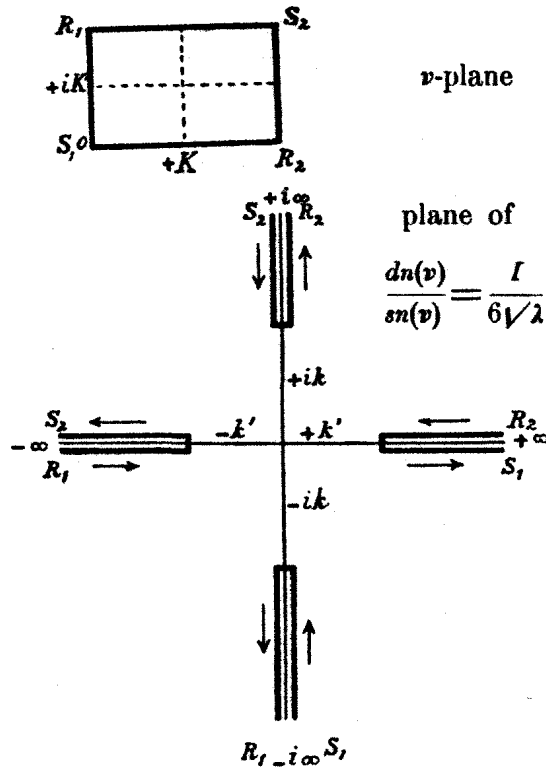


Fig. 1

$I_4 = 6 \sqrt{\frac{1-\lambda}{2}} = 6i \sqrt{\frac{\lambda-1}{2}}$ ; the quotient  $\frac{dn(v)}{sn(v)}$  assumes in those points successively the values  $0, \infty, \pm k', \pm ik$ . The corresponding values of  $v$  are congruent (*mod.*  $2K$  and  $2iK'$ ) with  $K + iK', 0, K$  and  $iK'$ . (see fig. 1).

*Case IVa.* Point  $O$  lies in the domain between the hyperbola and the asymptotes. The points  $R_1, R_2, S_1, S_2, T_1$  and  $T_2$  are all real;  $T_1$  and  $T_2$  lie both on the same side of the polar line of  $O$  as  $O$  itself. We shall assume that the polar line intersects that branch on which  $T_1$  lies. The order of the singular points is then  $S_1, R_1, T_1, R_2, S_2, T_2, S_1$ .

The values  $I_1 = 0, I_2 = \infty, I_3 = 6 \sqrt{\frac{1+\lambda}{2}}, I_4 = 6 \sqrt{\frac{1-\lambda}{2}}$  correspond resp. to the values of  $0, \infty, \pm 1$  and  $\pm k$  for  $\frac{1}{sn(v)}$ , thus to the values of  $v$  which are congruent (*mod.*  $2K$  and  $2iK'$ ) resp. with  $iK', 0, K$  and  $K + iK'$ . (see fig. 2).

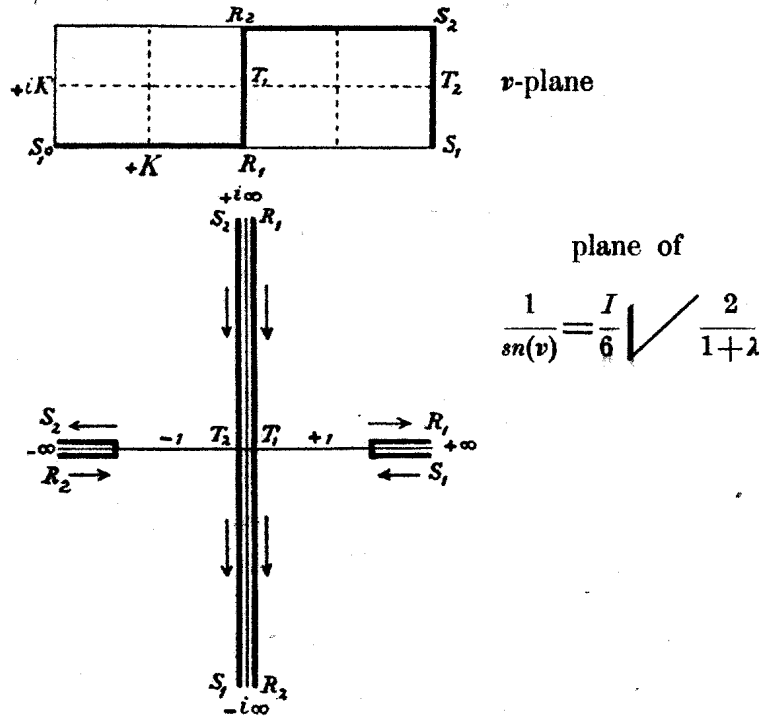


Fig. 2.

Case IVb. Point  $O$  lies inside the ellipse;  $T_1$  and  $T_2$  are real,  $R_1, R_2, S_1$  and  $S_2$  are imaginary.

IVb	in $T_1$	on $T_1 T_2$	in $T_2$	on $T_2 T_1$	in $T_1$
$v$	$iK'$	$iK' + \text{real}$	$2K + iK'$	$iK' + \text{real}$	$4K + iK'$
$\xi$	$\frac{\sqrt{1+\lambda} + \sqrt{1-\lambda}}{\lambda\sqrt{2}}$	pos. real	$\frac{\sqrt{1+\lambda} - \sqrt{1-\lambda}}{\lambda\sqrt{2}}$	pos. real	$\frac{\sqrt{1+\lambda} + \sqrt{1-\lambda}}{\lambda\sqrt{2}}$
$\zeta$	0	neg. real	0	pos. real	0
$I$	0	pos. real	0	neg. real	0

The points where the motion of  $I$  changes its sign are according to what was found in IVa the points for which  $\frac{1}{sn(v)} = \pm k$ , thus  $v = K + iK' \pmod{2K \text{ and } 2iK'}$  (see fig. 3).

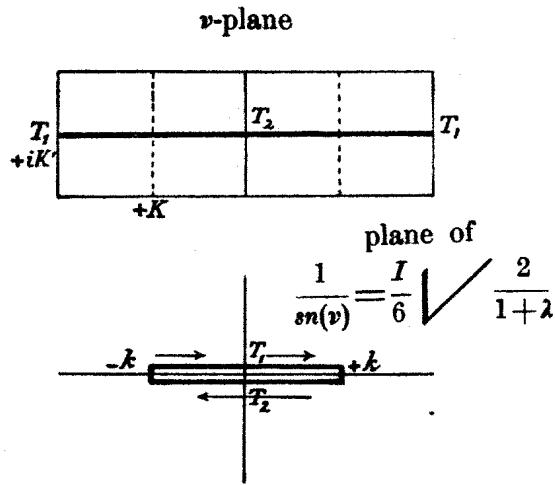


Fig. 3.

*Case VIa.* Point  $O$  lies on the concave side of the hyperbola:  $S_1, S_2, T_1$  and  $T_2$  are real,  $R_1$  and  $R_2$  are imaginary. Let  $T_1$  be the point of intersection of the diameter through  $O$  lying on the same side of the polar line as  $O$  itself.

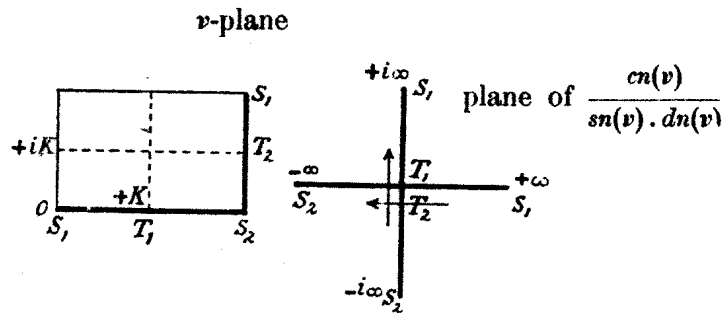


Fig. 4.

The values  $I_1=0, I_2=\infty, I_3=6 \sqrt{\frac{1+i\lambda'}{2}}, I_4=6 \sqrt{\frac{1-i\lambda'}{2}}$  correspond here respectively to the values  $0, \infty, \pm \sqrt[4]{\frac{1+i\lambda'}{1-i\lambda'}}$  and  $\pm \sqrt[4]{\frac{1-i\lambda'}{1+i\lambda'}}$  for  $\frac{cn(v)}{sn(v) \cdot dn(v)}$ , thus to the values of  $v$  which are congruent (mod.  $2K$  and  $K+iK'$ ) with  $K, 0, \frac{1}{2}(K+3iK'), \frac{1}{2}(K+iK')$  (see fig. 4).

*Case VIb.* Point  $O$  lies outside the ellipse;  $R_1, R_2, T_1$  and  $T_2$  are real,  $S_1$  and  $S_2$  are imaginary. The point of intersection  $T_1$  may lie on the same side of the polar line as  $O$  itself.

For the particular values of  $I$  and the corresponding values of  $v$  we can refer to VIa. (see fig. 5).

$V/a$	in $S_1^\infty$	on $S_1 T_1$	in $T_1$	on $T_1 S_2^\infty$	in $S_2^\infty$	on $S_2^\infty T_2$	in $T_2$	on $T_2 S_1^\infty$	in $S_1^\infty$
$\infty$	0	real	$K + \frac{\sqrt{-1 + \sqrt{1 + \lambda^2}}}{\lambda}$	real	$2K$	$2K + p.\text{imag.}$	$2K + iK' + \frac{\sqrt{+1 + \sqrt{1 + \lambda^2}}}{\lambda}$	$2K + p.\text{imag.}$	$2K + 2iK'$
$\infty$	$\infty$	pos. real	0	pos. real	$\infty$	pos. imag.	0	pos. imag.	$\infty$
$\infty$	$\infty$	neg. real	0	pos. real	$\infty$	neg. real	0	pos. real	$\infty$
$I$	$\infty$	pos. real	0	neg. real	$\infty$	neg. imag.	0	pos. imag.	$\infty$

$V/b$	in $R_1$	on $R_1 T_1$	in $T_1$	on $T_1 R_2$	in $R_2$	on $R_2 T_2$	in $T_2$	on $T_2 R_1$	in $R_1$
$\infty$	$K + iK'$	$K + p.\text{imag.}$	$K + 2iK' - \frac{\sqrt{-1 + \sqrt{1 + \lambda^2}}}{\lambda}$	$K + p.\text{imag.}$	$K + 3iK'$	$3iK' + \text{real}$	$2K + 3iK' - \frac{\sqrt{+1 + \sqrt{1 + \lambda^2}}}{\lambda}$	$3iK' + \text{real}$	$3K + 3iK'$
$\infty$	0	neg. real	0	neg. real	0	neg. imag.	0	neg. imag.	0
$\infty$	$+\frac{i}{2}$	pos. imag.	0	neg. imag.	$-\frac{i}{2}$	neg. imag.	0	pos. imag.	$+\frac{i}{2}$
$I$	$\infty$	pos. imag.	0	neg. imag.	$\infty$	neg. real	0	pos. real.	$\infty$

Before investigating the cases of degeneration III and V we shall occupy ourselves for a moment with the relation (53) (4<sup>th</sup> comm. p. 1011), existing between  $I$  and  $I^2$ . In the case of the conic it takes the shape of (65) (4<sup>th</sup> comm. p. 1018). The curve it represents is as can be expected symmetrical with respect to the  $X$ -axis ( $X = I^2$ ). To simplify the reasoning we shall translate the curve  $\Phi(X, Y) \equiv \Phi(I^2, I) = 0$  parallel to the  $X$ -axis and we shall decrease it and that by the formulae of transformation

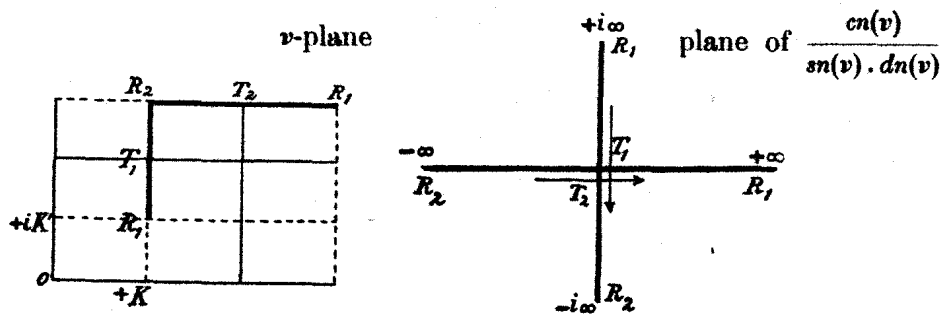


Fig. 5.

$$I^2 - 18 = 36\xi,$$

$$I = 6\eta.$$

The equation of the curve transformed in this manner runs as follows:

$$\Phi(\xi, \eta) \equiv \xi^2 - \eta^2 - \frac{\lambda^2}{4} = 0;$$

the curve is therefore a rectangular hyperbola. In the cases II and IV the  $\xi$ -axis is the real axis, in case VI the  $\eta$ -axis is the real axis. Each point of the conic  $F(x, y) = 0$  corresponds to one point of this rectangular hyperbola whilst to one point of  $\Phi = 0$  two points of  $F = 0$  are conjugated. The points for which  $I = 0$  have as absciss  $\xi = -\frac{1}{2}$ . The line  $\xi = -\frac{1}{2}$  does not intersect the curve  $\Phi$  in case II, but it does in the cases IV and VI. The point at infinity on  $\xi + \eta = 0$  represents the points  $S_1$  and  $S_2$ ; the point at infinity on  $\xi - \eta = 0$  represents the two points  $R_1$  and  $R_2$ . The points  $T_1$  and  $T_2$  are represented by the points of intersection of  $\Phi = 0$  with  $\xi = -\frac{1}{2}$ . The images of the points  $T_1$  and  $T_2$  are in case VI united in the point of intersection of  $\xi = -\frac{1}{2}$  with the branch of  $\Phi = 0$  lying under the  $\xi$ -axis. The images of  $T_1$  and  $T_2$  are always points where the motion changes its sign along the curve  $\Phi$ .

Now we have to investigate the cases of degeneration.

Case III<sup>a</sup>.  $\lambda = +1$ ,  $\sigma_1 = 0$ ,  $a_1$ , and  $a_2$ , not disappearing at the same time.

The point  $O$  lies on one of the asymptotes, without coinciding with the centre. So this position occurs with the hyperbola only.

Here equation (71) holds, in which is put  $\tau_0 = 0$ ,

$$I = \pm \frac{6}{\sin \tau} \dots \dots \dots (71)$$

Equation (62) (4<sup>th</sup> comm. p. 1015) passes, on account of the relation

$$a_{22}A_{22} = \Delta,$$

and with the aid of (72), into

$$I^2 = \frac{9}{\zeta^2} (-\zeta^4 + 2\zeta^2 - 1) = -\frac{9(\zeta^2 - 1)^2}{\zeta^2},$$

from which ensues, in connection with (71),

$$\frac{6}{\sin \tau} = \pm \frac{3i(\zeta^2 - 1)}{\zeta},$$

or

$$\zeta = \pm \frac{i(1 \pm \cos \tau)}{\sin \tau}.$$

We choose for  $\zeta$ :

$$\zeta = +i \frac{1 - \cos \tau}{\sin \tau} = +i \operatorname{tg} \frac{\tau}{2}, \dots \dots \dots (85)$$

and find in this manner

$$\zeta = +\frac{i}{2} \sec^2 \frac{\tau}{2}.$$

Now the equations (76) and (77) are incompatible. If they depended on each other we should have  $A_{22} = 0$ , which has not been supposed to be the case.

Equation (77) now runs:

$$a_{11}x + a_{22}y = -a_{22}z \sec^2 \frac{\tau}{2} \dots \dots \dots (86)$$

Bringing this equation into connection with  $F(x, y, z) = 0$ , we find

$$\begin{aligned} 2A_{11} \sec^2 \frac{\tau}{2} \cdot x &= \left\{ a_{22}a_{22} \sec^4 \frac{\tau}{2} - a_{22}^2 \left( 2 \sec^2 \frac{\tau}{2} - 1 \right) \right\} z \\ 2A_{22} \sec^2 \frac{\tau}{2} \cdot y &= \left\{ a_{11}a_{22} \sec^4 \frac{\tau}{2} - a_{11}^2 \left( 2 \sec^2 \frac{\tau}{2} - 1 \right) \right\} z \end{aligned} \dots \dots \dots (87)$$

These formulae can be used unless either  $A_{11}$  or  $A_{22}$  is zero. Therefore we will mention also the expressions for  $x$  and  $y$  for the case  $A_{11} = 0$ . Then we have  $a_{22} = 0$  on account of  $a_{11}A_{11} + a_{22}A_{22} = 0$ . We then find immediately out of (86) the expression for  $x$ , out of the second equation (87) in which  $A_{22}$  is replaced by  $a_{11}a_{11}$  the expression for  $y$ . So the solution is:

$$\begin{aligned} a_{11}x &= -a_{22}z \cdot \sec^2 \frac{\tau}{2} \\ 2a_{11}a_{11} \sec^2 \frac{\tau}{2} y &= \left\{ a_{11}a_{22} \sec^4 \frac{\tau}{2} - a_{11}^2 \left( 2 \sec^2 \frac{\tau}{2} - 1 \right) \right\} z. \end{aligned}$$

Case III<sup>b</sup>.  $a_{11} = a_{22} = 0$ .

The point  $O$  coincides with the centre.

Now we have

$$I = 0.$$

The expressions for  $x$  and  $y$  are of the form:

$$x = (\alpha e^{i\tau} + \alpha' e^{-i\tau}) z,$$

$$y = (\beta e^{i\tau} + \beta' e^{-i\tau}) z.$$

In order to have  $F \equiv a_{11}x^2 + 2a_{12}xy + a_{22}y^2 + a_{33}z^2 = 0$ , we must put:

$$\alpha = \sigma(-a_{12} + \sqrt{-A_{33}}), \quad \alpha' = \sigma'(-a_{12} - \sqrt{-A_{33}}),$$

$$\beta = \sigma a_{11}, \quad \beta' = \sigma' a_{11}$$

with the condition

$$\sigma\sigma' = \frac{a_{33}}{-4a_{11}A_{33}}.$$

In the case of the real ellipse we have  $A_{33} > 0$  and  $\frac{a_{33}}{a_{11}} < 0$ . We then can put:

$$\sigma = \sigma' = \frac{1}{2} \sqrt{\frac{-a_{33}}{a_{11}A_{33}}}.$$

So we find

$$\left. \begin{aligned} x &= \frac{1}{2} \sqrt{\frac{-a_{33}}{a_{11}A_{33}}} \cdot \{-a_{12}(e^{i\tau} + e^{-i\tau}) + i\sqrt{A_{33}}(e^{i\tau} - e^{-i\tau})\} z = \\ &= \sqrt{\frac{-a_{33}}{a_{11}A_{33}}} \cdot (-a_{12} \cos \tau - \sqrt{A_{33}} \sin \tau) z, \\ y &= \frac{1}{2} \sqrt{\frac{-a_{33}}{a_{11}A_{33}}} \cdot a_{11}(e^{i\tau} + e^{-i\tau}) z = \sqrt{\frac{-a_{11}a_{33}}{A_{33}}} \cdot \cos \tau \cdot z. \end{aligned} \right\} (88)$$

We can use the same expression if we have to deal with a hyperbola not intersecting the  $x$ -axis. For then  $A_{33} < 0$  and  $\frac{a_{33}}{a_{11}} > 0$ , so  $\sigma = \sigma'$  real. We prefer to write  $-\sqrt{-A_{33}} \cdot sh(i\tau)$  for  $\sqrt{A_{33}} \cdot \sin \tau = -i\sqrt{-A_{33}} \cdot \sin \tau$  and  $ch(i\tau) \cos$  for  $\tau$ . Then real points of the hyperbola correspond to purely imaginary values of  $\tau$ .

If the hyperbola does intersect the  $x$ -axis we have  $A_{33} < 0$  and  $\frac{a_{33}}{a_{11}} < 0$ , so  $\sigma = \sigma'$  imaginary.

We then put  $\sigma = -\sigma' = \frac{1}{2} \sqrt{\frac{a_{33}}{a_{11}A_{33}}}$  and get in this manner

$$\left. \begin{aligned} x &= \frac{1}{2} \sqrt{\frac{a_{33}}{a_{11}A_{33}}} \cdot \{-a_{12}(e^{i\tau} - e^{-i\tau}) + \sqrt{-A_{33}}(e^{i\tau} + e^{-i\tau})\} z = \\ &= \sqrt{\frac{a_{33}}{a_{11}A_{33}}} \cdot \{-a_{12} sh(i\tau) + \sqrt{A_{33}} ch(i\tau)\} z, \\ y &= \frac{1}{2} \sqrt{\frac{a_{33}}{a_{11}A_{33}}} \cdot a_{11}(e^{i\tau} - e^{-i\tau}) z = \sqrt{\frac{a_{11}a_{33}}{A_{33}}} \cdot sh(i\tau) \cdot z. \end{aligned} \right\} (88)'$$



Here also  $\tau$  must describe in its complex plane the imaginary axis.

For  $a_{11} = 0$ , we get  $(2a_{12}x + a_{22}y)y + a_{33} = 0$ .

A solution of this is given by

$$x = \frac{-1}{2a_{12}} (a_{22}e^{i\tau} + a_{33}e^{-i\tau}),$$

$$y = e^{i\tau}.$$

Here also only purely imaginary values of  $\tau$  come in consideration, as might be expected.

The second case of degeneration (IV) presents itself for  $\lambda = 0$ , i. e.  $d_1 = +1$ . Here we must distinguish three subdivisational cases, viz.

*IV<sup>a</sup>.*  $a_{33} = 0$ : the point  $O$  lies on the conic,

*IV<sup>b</sup>.*  $A_{33} = 0$ : the conic is a parabola,

*IV<sup>c</sup>.*  $a_{33} = 0$  and  $A_{33} = 0$ : the point  $O$  lies on the parabola.

*Case IV<sup>a</sup>.* Here we have (70a) (4<sup>th</sup> comm. p. 1017); substitution of  $\tau_0 = 0$  furnishes

$$I = + 3\sqrt{2} \cdot th \frac{\tau}{\sqrt{2}}, \dots \dots \dots (70'a)$$

so

$$i = \frac{3}{ch^2 \frac{\tau}{\sqrt{2}}}$$

Now the equations (62) and (63) (4<sup>th</sup> comm. p. 1015) teach us

$$g = a_{11}x + a_{21}y = \frac{2\Delta \dot{I}z}{3A_{33}} = \frac{2\Delta z}{A_{33}} \cdot \frac{1}{ch^2 \frac{\tau}{\sqrt{2}}}$$

$$A_{22}x - A_{12}y = \sqrt{-A_{33}g^2 + 2\Delta gz} = \sqrt{\left( -\frac{4\Delta^2}{A_{33}^2} \cdot \frac{1}{ch^4 \frac{\tau}{\sqrt{2}}} + \frac{4\Delta^2}{A_{33}^2} \cdot \frac{1}{ch^2 \frac{\tau}{\sqrt{2}}} \right) \cdot z}$$

$$= \frac{2\Delta z}{A_{33}} \cdot \frac{sh \frac{\tau}{\sqrt{2}}}{ch^2 \frac{\tau}{\sqrt{2}}};$$

so we get

$$\left. \begin{aligned} x &= \frac{2z}{A_{33}ch^2 \frac{\tau}{\sqrt{2}}} \left( A_{11} + a_{33} \sqrt{A_{33}} \cdot sh \frac{\tau}{\sqrt{2}} \right) \\ y &= \frac{2z}{A_{33}ch^2 \frac{\tau}{\sqrt{2}}} \left( A_{22} - a_{12} \sqrt{A_{33}} \cdot sh \frac{\tau}{\sqrt{2}} \right) \end{aligned} \right\} \dots \dots (89)$$

In the case  $A_{11} < 0$  we prefer to write  $i\sqrt{-A_{11}} \cdot sh \frac{\tau}{\sqrt{2}} =$   
 $= +\sqrt{-A_{11}} \cdot sin \frac{i\tau}{\sqrt{2}}$  for  $\sqrt{A_{11}} \cdot sh \frac{\tau}{\sqrt{2}}$  and  $cos \frac{i\tau}{\sqrt{2}}$  for  $ch \frac{\tau}{\sqrt{2}}$ .

So, whilst the formulae (89) are specially suitable for the ellipse we do better in using for the hyperbola

$$\left. \begin{aligned} x &= \frac{2z}{A_{11} \cos^2 \frac{i\tau}{\sqrt{2}}} \left( A_{11} + a_{11} \sqrt{-A_{11}} \cdot sin \frac{i\tau}{\sqrt{2}} \right) \\ y &= \frac{2z}{A_{11} \cos^2 \frac{i\tau}{\sqrt{2}}} \left( A_{11} - a_{11} \sqrt{-A_{11}} \cdot sin \frac{i\tau}{\sqrt{2}} \right) \end{aligned} \right\} \dots (89')$$

Consequently the real points of the hyperbola correspond to purely imaginary values of  $\tau$ .

Case IV<sup>b</sup>. Putting  $\tau_0 = 0$ , (70b) (4<sup>th</sup> comm. p. 1017) we find

$$I = -3\sqrt{2} \cdot th \frac{\tau}{\sqrt{2}} \dots (70'b)$$

and therefore

$$i = -\frac{3}{ch^2 \frac{\tau}{\sqrt{2}}}$$

So the formulae (62) and (63) now give

$$g = a_{11}x + a_{22}y + a_{33}z = -\frac{3a_{33}z}{2i} = \frac{a_{33}z}{2} ch^2 \frac{\tau}{\sqrt{2}},$$

i.e.

$$a_{11}x + a_{22}y = \frac{a_{33}z}{2} \left( ch^2 \frac{\tau}{\sqrt{2}} - 2 \right)$$

and

$$A_{22}x - A_{11}y = \sqrt{2\Delta gz - \Delta a_{33}z^2} = \sqrt{a_{33}\Delta} \cdot z \cdot sh \frac{\tau}{\sqrt{2}},$$

so we find

$$\left. \begin{aligned} x &= \left\{ \frac{a_{33} \sqrt{a_{33}\Delta}}{\Delta} sh \frac{\tau}{\sqrt{2}} + \frac{a_{33}A_{11}}{2\Delta} \left( ch^2 \frac{\tau}{\sqrt{2}} - 2 \right) \right\} z, \\ y &= \left\{ \frac{-a_{11} \sqrt{a_{33}\Delta}}{\Delta} sh \frac{\tau}{\sqrt{2}} + \frac{a_{33}A_{22}}{2\Delta} \left( ch^2 \frac{\tau}{\sqrt{2}} - 2 \right) \right\} z. \end{aligned} \right\} \dots (90)$$

Case IV<sup>c</sup>. Here we have

$$I = \pm 3\sqrt{2}$$

The equation

$$F \equiv a_{11}x^2 + 2a_{12}xy + \frac{a_{12}^2}{a_{11}}y^2 + 2a_{13}x + 2a_{23}y = 0$$

or

$$(a_{11}x + a_{12}y)^2 + 2a_{11}(a_{13}x + a_{23}y) = 0$$

passes by the substitution

$$\left. \begin{aligned} a_{11}x + a_{12}y &= 2a_{11}\xi \\ a_{13}x + a_{23}y &= -2a_{11}\eta \end{aligned} \right\} \dots \dots \dots (91)$$

into

$$\eta = \xi^2,$$

a solution of which (see 2<sup>nd</sup> comm. p. 590) is

$$\xi = e^{-\frac{\tau}{\sqrt{2}}}, \quad \eta = e^{-\tau\sqrt{2}} \dots \dots \dots (92)$$

Out of (91) and (92) we deduce

$$\left. \begin{aligned} x &= -\frac{2a_{11}}{A_{12}} \left( a_{23} e^{-\frac{\tau}{\sqrt{2}}} + a_{13} e^{-\tau\sqrt{2}} \right) \\ y &= \frac{2a_{11}}{A_{12}} \left( a_{13} e^{-\frac{\tau}{\sqrt{2}}} + a_{11} e^{-\tau\sqrt{2}} \right) \end{aligned} \right\} \dots \dots \dots (93)$$

These formulae are always applicable, as the supposition  $A_{12} = 0$  would imply the degeneration of the parabola.

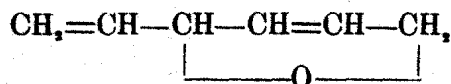
**Chemistry.** — “*On some internal unsaturated ethers*”. By J. W. LE HEUX. (communicated by Prof. VAN ROMBURGH).

(Preliminary communication).

(Communicated in the meeting of April 26, 1912).

By the action of formic acid on mannitol FAUCONNIER obtained a mixture of formic esters of this hexavalent alcohol, which submitted to dry distillation, yielded among other products a liquid of the composition  $C_6H_8O$ , boiling at  $107^\circ$ — $109^\circ$ .

VAN MAANEN (Dissertation, Utrecht 1909) who investigated this substance and mentions it as a liquid boiling at  $107^\circ$  proposed as the most probable structural formula:



As the mode of formation of this substance does not give a complete insight into its structural formula, Prof. VAN ROMBURGH proposed to me to prepare the various possible oxides of hexadiene by other