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Mathematics. — Prof. KLUYVER presents a paper: “*Evaluation of two definite integrals.*”

Supposing x to be real the integrals

$$f(x, m) = \int_0^{\infty} \frac{\cos xt}{(1+t^2)^m} dt \quad \text{and} \quad \varphi(x, m) = \int_0^{\infty} \frac{\sin xt}{(1+t^2)^m} dt$$

will have a definite meaning, if only the real part of the parameter m be positive. In what follows we will show how to expand these integrals into rapidly converging power series.

The first of them, the integral $f(x, m)$, is a particular solution of the linear equation

$$x \frac{d^2 y}{dx^2} - 2(m-1) \frac{dy}{dx} - xy = 0,$$

the primitive of which, involving two arbitrary constants A and B , may be written in the form

$$y = A L(x, m) + B x^{2m-1} M(x, m),$$

where

$$L(x, m) = \sum_{h=0}^{h=\infty} \frac{\left(\frac{x}{2}\right)^{2h}}{h! \Gamma(-m + \frac{3}{2} + h)},$$

$$M(x, m) = \sum_{h=0}^{h=\infty} \frac{\left(\frac{x}{2}\right)^{2h}}{h! \Gamma(m + \frac{1}{2} + h)},$$

and the constants A and B must now be determined so that y represents the function $f(x, m)$. To find A , we suppose $m > \frac{1}{2}$ and put x equal to zero. In that case we have

$$f(0, m) = \int_0^{\infty} \frac{dt}{(1+t^2)^m} = \frac{\Gamma(\frac{1}{2}) \Gamma(m - \frac{1}{2})}{2 \Gamma(m)} = A L(0, m) = \frac{A}{\Gamma(-m + \frac{1}{2})}$$

and hence

$$A = - \frac{\pi}{2 \cos \pi m} \cdot \frac{\Gamma(\frac{1}{2})}{\Gamma(m)}.$$

For the deduction of the constant B it is convenient to consider first the function $f'(x, m)$ in another form. Let the real part of m still be positive, then we have

$$\frac{\Gamma(m)}{(1+t^2)^n} = \int_0^\infty e^{-u(1+t^2)} u^{m-1} du,$$

and hence

$$\Gamma(m) f(x, m) = \int_0^\infty e^{-u} u^{m-1} du \int_0^\infty e^{-ut^2} \cos xt dt = \frac{1}{2} \sqrt{\pi} \int_0^\infty e^{-u - \frac{x^2}{4u}} u^{m-\frac{3}{2}} du.$$

From the latter integral a simple functional relation is derivable. Changing the variable u into $\frac{x^2}{4v}$ we may write

$$\begin{aligned} \Gamma(m) f(x, m) &= \frac{1}{2} \sqrt{\pi} \left(\frac{x}{2}\right)^{2m-1} \int_0^\infty e^{-v - \frac{x^2}{4v}} v^{-m-\frac{1}{2}} dv = \\ &= \left(\frac{x}{2}\right)^{2m-1} \Gamma(1-m) f(x, 1-m) \end{aligned}$$

and so it follows that the function

$$\begin{aligned} \frac{\Gamma(m) f(x, m)}{\left(\frac{x}{2}\right)^m} &= -\frac{\pi}{2 \cos \pi m} \Gamma\left(\frac{1}{2}\right) \left(\frac{x}{2}\right)^{-m} L(x, m) + \\ &+ 2^{2m-1} B \Gamma(m) \left(\frac{x}{2}\right)^{m-1} M(x, m) \end{aligned}$$

remains unaltered, if m is replaced by $1-m$.

Now obviously the series L and M are connected by the relation

$$L(x, 1-m) = M(x, m),$$

hence we must have

$$\begin{aligned} 2^{2m-1} B \Gamma(m) &= + \frac{\pi}{2 \cos \pi m} \Gamma\left(\frac{1}{2}\right), \\ B &= \frac{\pi}{2 \cos \pi m} \cdot \frac{\Gamma\left(\frac{1}{2}\right)}{\Gamma(m)} \left(\frac{1}{2}\right)^{2m-1} \end{aligned}$$

and therefore

$$f(x, m) = \int_0^\infty \frac{\cos xt}{(1+t^2)^m} dt = \frac{\pi}{2 \cos \pi m} \cdot \frac{\Gamma\left(\frac{1}{2}\right)}{\Gamma(m)} \left\{ -L(x, m) + \left(\frac{x}{2}\right)^{2m-1} M(x, m) \right\}.$$

Now it will be observed, that the series $L(x, m)$ and $M(x, m)$ converge for all values of x and m , and so we must conclude, that the function $f(x, m)$ exists over the whole x -plane, that its only singularities are $x=0$ and $x=\infty$, and that therefore the integral, we started with, represents the function in a very incomplete manner.

Numerical evaluation of the integral for not too large values of x offers no difficulties, as the series $L(x, m)$ and $M(x, m)$ converge rapidly. Because of the equation

$$\Gamma(m) f(x, m) = \frac{1}{2} \sqrt{\pi} \int_0^{\infty} e^{-u - \frac{x^2}{4u}} u^{m - \frac{3}{2}} du$$

the result will always be a positive number and the integral will not vanish for any real value of x .

A few further remarks may be made. Firstly we may state, that $f(x, m)$ is intimately connected with BESSEL'S function $J_n(x)$. In fact, by means of the usual expansion of $J_n(x)$ we may verify the relation

$$f(x, m) = - \frac{\pi}{2 \cos \pi m} \cdot \frac{\Gamma(\frac{1}{2})}{\Gamma(m)} \left(\frac{x}{2}\right)^{m - \frac{1}{2}} \left\{ e^{-\frac{\pi i}{2}(m - \frac{1}{2})} J_{-m + \frac{1}{2}} \left(xe^{-\frac{\pi i}{2}}\right) - e^{+\frac{\pi i}{2}(m - \frac{1}{2})} J_{m - \frac{1}{2}} \left(xe^{-\frac{\pi i}{2}}\right) \right\}.$$

From this we infer, that for positive integer values of m the origin $x = 0$ ceases to be a singular point, and that $f(x, m)$ can be expressed in finite terms. We shall find by actual substitution of the finite

expressions for $J_{-m + \frac{1}{2}} \left(xe^{-\frac{\pi i}{2}}\right)$ and $J_{m - \frac{1}{2}} \left(xe^{-\frac{\pi i}{2}}\right)$

$$f(x, m) = \frac{\pi}{2} \cdot \frac{e^{-x}}{(m-1)!} \left(\frac{x}{2}\right)^{m-1} \sum_{h=0}^{h=m-1} \frac{(m-1+h)!}{h!(m-1-h)!} \cdot \left(\frac{1}{2x}\right)^h.$$

However this result may be obtained in a simpler way as follows. It can be shewn, that $f(x, m)$ obeys the relation

$$D_{x=1}^h \left\{ x^{-m + \frac{1}{2}} f(x \sqrt{\alpha}, m) \right\} = \frac{(-1)^h \Gamma(m+h)}{\Gamma(m)} f(x, m+h),$$

and since we have

$$f(x \sqrt{\alpha}, 1) = \int \frac{\cos \alpha t \sqrt{\alpha}}{1+t^2} dt = \frac{\pi}{2} e^{-x \sqrt{\alpha}},$$

we get for all positive values of m

$$f(x, m) = \frac{\pi}{2} \cdot \frac{(-1)^{m-1}}{(m-1)!} D_{\sigma=1}^{m-1} \left\{ \frac{e^{-x \sqrt{\alpha}}}{\sqrt{\alpha}} \right\},$$

a result that can be identified with that obtained before.

The singularity in the origin $x = 0$ becomes logarithmic, when $2m$ is an odd integer $2k + 1$. The expression of $f(x, m)$ is in this

case somewhat intricate. By repeated differentiations it is derived from $f\left(x, \frac{1}{2}\right)$, for we have

$$D_{x=1}^k f\left(x \vee \alpha, \frac{1}{2}\right) = \frac{(-1)^k \Gamma(k + \frac{1}{2})}{\Gamma(\frac{1}{2})} f\left(x, k + \frac{1}{2}\right).$$

To evaluate $f\left(x, \frac{1}{2}\right)$, we put $m = \frac{1}{2} + \sigma$ in the general expression for $f(x, m)$, and make σ tend to zero. In this way we get

$$f\left(x, \frac{1}{2}\right) = \lim_{\sigma \rightarrow 0} \frac{\pi}{2 \sin \pi \sigma} \left\{ \sum_{h=0}^{\infty} \frac{\left(\frac{x}{2}\right)^{2h}}{h! \Gamma(h+1-\sigma)} - \left(\frac{x}{2}\right)^{2\sigma} \sum_{h=0}^{\infty} \frac{\left(\frac{x}{2}\right)^{2h}}{h! \Gamma(h+1+\sigma)} \right\},$$

$$f\left(x, \frac{1}{2}\right) = \int_0^{\infty} \frac{\cos xt}{\sqrt{1+t^2}} dt = \sum_{h=0}^{\infty} \frac{\left(\frac{x}{2}\right)^{2h}}{(h!)^2} \left(\frac{1}{1} + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{h} - C - \log \frac{x}{2} \right).$$

We shall now pass on to consider the second integral

$$\varphi(x, m) = \int_0^{\infty} \frac{\sin xt}{(1+t^2)^m} dt,$$

and it will appear that its character is quite similar to that of the first. Again we transform the integral by the aid of the identity

$$\frac{\Gamma(m)}{(1+t^2)^m} = \int_0^{\infty} e^{-u(1+t^2)} u^{m-1} du$$

and obtain

$$\Gamma(m) \varphi(x, m) = \int_0^{\infty} e^{-u} u^{m-1} du \int_0^{\infty} e^{-ut^2} \sin xt dt.$$

A further transformation gives

$$\int_0^{\infty} e^{-ut^2} \sin xt dt = \frac{x}{2u} \int_0^1 e^{-\frac{x^2}{4u}(1-t^2)} dt = \frac{x}{4u} \int_0^1 e^{-\frac{x^2 w}{4u}} (1-w)^{-\frac{1}{2}} dw.$$

and therefore

$$\Gamma(m) \varphi(x, m) = \frac{x}{4} \int_0^1 (1-w)^{-\frac{1}{2}} dw \int_0^{\infty} e^{-u - \frac{x^2 w}{4u}} u^{m-2} du.$$

Comparing this equation with the equality obtained before

$$\Gamma(m) f(x, m) = \frac{1}{2} \sqrt{\pi} \int_0^{\infty} e^{-u - \frac{x^2}{4u}} u^{m - \frac{3}{2}} du,$$

it follows, that we may write

$$\Gamma(m) \varphi(x, m) = \frac{x}{2\sqrt{\pi}} \int_0^1 (1-w)^{-\frac{1}{2}} dw \Gamma(m - \frac{1}{2}) f\left(x\sqrt{w}, m - \frac{1}{2}\right)^1).$$

We now expand $\Gamma(m - \frac{1}{2}) f\left(x\sqrt{w}, m - \frac{1}{2}\right)$ and make the substitution

$$\Gamma(m - \frac{1}{2}) f\left(x\sqrt{w}, m - \frac{1}{2}\right) = \frac{\pi}{2\sin \pi m} \cdot \Gamma(\frac{1}{2}) \left\{ -L\left(x\sqrt{w}, m - \frac{1}{2}\right) + \left(\frac{x}{2}\right)^{2m-2} M\left(x\sqrt{w}, m - \frac{1}{2}\right) \right\}.$$

Then integrating with respect to w , we find the desired expansion of $\varphi(x, m)$ in the form

$$\varphi(x, m) = \int_0^{\infty} \frac{\sin xt}{(1+t^2)^m} dt = \frac{\pi}{2\sin \pi m} \cdot \frac{\Gamma(\frac{1}{2})}{\Gamma(m)} \left\{ -N(x, m) + \left(\frac{x}{2}\right)^{2m-1} M(x, m) \right\},$$

where $N(x, m)$ represents the new series

$$N(x, m) = \sum_{h=0}^{h=\infty} \frac{\left(\frac{x}{2}\right)^{2h+1}}{\Gamma(h + \frac{1}{2}) \Gamma(-m + 2 + h)}.$$

The same remarks as were made concerning the first integral $f(x, m)$, can here be made again. The integral has only a meaning for real values of x and for positive values of m , but from the expansion is inferred, that the integral incompletely represents a function of x which exists over the whole x -plane, quite independently of the values assigned to the parameter m . Again the origin $x=0$ and $x=\infty$ are the only singularities of the function. The singularities are logarithmic, when m is an integer and the origin becomes a regular point, when $2m$ is equal to an odd integer

¹⁾ It is possible to invert this relation. It may be shewn that we have also

$$\begin{aligned} \Gamma(m) f(x, m) - \frac{1}{2} \sqrt{\pi} \Gamma(m - \frac{1}{2}) &= \\ = -\frac{x}{2\sqrt{\pi}} \int_0^1 (1-w)^{-\frac{1}{2}} dw \Gamma(m - \frac{1}{2}) \varphi\left(x\sqrt{w}, m - \frac{1}{2}\right). \end{aligned}$$

$2k+1$, but in no case is a finite expression by means of elementary functions obtainable.

The function $\varphi(x, m)$ as well as $f'(x, m)$ satisfies the relation

$$D_{x=1}^h \left\{ x^{-m+\frac{1}{2}} \varphi(x \sqrt{\alpha}, m) \right\} = \frac{(-1)^h \Gamma(m+h)}{\Gamma(m)} \varphi(x, m+h),$$

and by means of this rule expansions for $\varphi(x, k)$ and $\varphi\left(x, k + \frac{1}{2}\right)$ may be deduced from the equations

$$\begin{aligned} \varphi(x, 1) &= \int_0^{\infty} \frac{\sin xt}{1+t^2} dt = \sum_{h=0}^{\infty} \frac{x^{2h+1}}{(2h+1)!} \left(\frac{1}{1} + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{2h+1} - C^{-\log x} \right) = \\ &= \frac{1}{2} \left\{ e^{-x} \text{Li}(e+x) - e+x \text{Li}(e-x) \right\}, \end{aligned}$$

and

$$\varphi\left(x, \frac{1}{2}\right) = \int_0^{\infty} \frac{\sin xt}{\sqrt{1+t^2}} dt = \frac{\pi}{2} \sum_{h=0}^{\infty} \frac{\left(-\frac{x}{2}\right)^h}{\Gamma\left(\frac{h}{2} + 1\right)^2}.$$

Botany. — On “*Leptostroma austriacum* OUD., a hitherto unknown *Leptostromacea* living on the needles of *Pinus austriaca*; and on *Hymenopsis Typhae* (Fuck.) SACC., a hitherto insufficiently described *Tuberculariacea*, occurring on the withered leafsheaths of *Typha latifolia*.” By Prof. C. A. J. A. OUDEMANS.

1. LEPTOSTROMA AUSTRIACUM OUD.

(Plate I.)

On the 13th of June 1904 I received from Dr. J. RITZEMA Bos, Professor at Amsterdam, a number of specimens of transplanted seedlings of *Pinus austriaca*, originating from Schoorl, all dead and of which the accompanying letter informed me that the roots showed here and there cushionlike prominences, the surface of which was covered with shuttle-shaped conidia, divided into cells, and the microscopic properties of which resembled most those of conidia of the genus *Fusarium*.

Besides I found, without my attention having been directed to it, that most needles of the dead plantlets were spotted on both