

**Mathematics.** — *Skew Correlation between Three and More Variables*, II. By Prof. M. J. VAN UVEN. (Communicated by Prof. A. A. NIJLAND).

(Communicated at the meeting of May 25, 1929).

We may still apply the method hitherto followed, when there is *no* linear correlation between the original variables  $x_1, x_2, x_3$  appearing in the given frequency distribution **A**, but when linear correlation can be found between the (unimodular) variables  $t_1, t_2, t_3$ ,  $t_a$  being a function of only  $x_a$  ( $a = 1, 2, 3$ ).

Considering, in particular,  $t_a \{x_a\}$  as an ever increasing one-valued function of  $x_a$  ( $a = 1, 2, 3$ )<sup>1)</sup>,  $s_a(k_a)$  (**B<sup>a</sup>**) gives the probability of  $x_a < x_a(k_a)$ . Thus the equation **B<sup>a</sup>** associates a value  $z_a(k_a)$  with  $x_a(k_a)$ , and we obtain a series of  $\nu_a - 1$  empirical values of the function  $z_a \{x_a\}$ .

Likewise  $\sigma_{b;a}(k_a - \frac{1}{2}, k_b)$  (**B(b; a)**) furnishes the probability of [ $x_b < x_b(k_b)$ , being given:  $x_a(k_a - 1) < x_a < x_a(k_a)$ ]. Hence the equation **B(b; a)** joins a value  $\zeta_{b;a}(k_a - \frac{1}{2}, k_b)$  to the set  $x_a(k_a - \frac{1}{2}), x_b(k_b)$ , and we obtain a series of empirical values of the function  $\zeta_{b;a} \{x_a, x_b\}$ .

As  $x_a$  itself is to be considered as a function of  $z_a$ , we may also conceive  $\zeta_{b;a}$  as a function of  $z_a$  and  $z_b$ .

Now  $z_a$  is proportional to  $t_a$ , so that we (cf 13) may put **9ter (b; a)** into the form:

$$\zeta_{b;a} = \frac{-A_{ba} \cdot \left(\frac{A_{aa}}{A}\right)^{1/2} \cdot z_a + A_{aa} \cdot \left(\frac{A_{bb}}{A}\right)^{1/2} \cdot z_b}{\sqrt{A_{aa}}} = \left. \begin{aligned} &= \frac{-A_{ba} \cdot z_a + \sqrt{A_{aa} A_{bb}} \cdot z_b}{\sqrt{A}} = -\cot \omega_c \cdot z_a + \frac{1}{\sin \omega_c} \cdot z_b \end{aligned} \right\} \text{9 quarter (b; a)}$$

Hence  $\zeta_{b;a}$  appears to be a *linear* function of  $z_a$  and  $z_b$ .

Finally  $S_c(k_a - \frac{1}{2}, k_b - \frac{1}{2}, k_c)$  (**B(c; ba)**) gives the probability of [ $x_c < x_c(k_c)$ , being given:  $x_a(k_a - 1) < x_a < x_a(k_a), x_b(k_b - 1) < x_b < x_b(k_b)$ ]. Thus the equation **B(c; ba)** associates a value  $Z_c(k_a - \frac{1}{2}, k_b - \frac{1}{2}, k_c)$  with the set  $x_a(k_a - \frac{1}{2}), x_b(k_b - \frac{1}{2}), x_c(k_c)$ . This furnishes us a set of empirical values of the function  $Z_c \{x_1, x_2, x_3\}$ . As  $x_a$  depends only on  $z_a$ ,  $Z_c$  is also to be considered as a function of all three variables  $z_1, z_2, z_3$ .

<sup>1)</sup> The limiting agreement that  $t \{x\}$  shall be monotone (increasing) is not necessary, but only desirable from a practical point of view and can in certain cases be dropped. (See: M. J. VAN UVEN, *Scheeve Frequentiekrommen*, Versl. K. A. v. W. 25, p. 709, *Skew Frequency Curves*, Proceed. K. Ak. v. Wet. Amsterdam, Vol. 19, p. 670).

In virtue of

$$Z_c = \lambda_{ca} \cdot t_a + \lambda_{bc} \cdot t_b + t_c = \frac{\lambda_{ca} \sqrt{A_{aa}} \cdot z_a + \lambda_{bc} \sqrt{A_{bb}} \cdot z_b + \sqrt{A_{cc}} \cdot z_c}{\sqrt{A}} \quad \left. \vphantom{\frac{\lambda_{ca} \sqrt{A_{aa}} \cdot z_a + \lambda_{bc} \sqrt{A_{bb}} \cdot z_b + \sqrt{A_{cc}} \cdot z_c}{\sqrt{A}}} \right\} \text{9 quater (c; ba)}$$

$Z_c$  is a linear function of  $z_1, z_2, z_3$ .

In constructing empirically the functions  $\zeta_{b;a}$  and  $Z_c$  it is necessary to determine the required mean values as exactly as possible.

Now the values of the functions  $\zeta_{b;a}, Z_c$ , which correspond to the sets of class-limits  $x_a(k_a)$ , and to the associated values  $z_a(k_a)$ , are to be calculated by interpolation. If these calculations are made correctly, and if the conditions that the unimodular variables  $t_1, t_2, t_3$  are pure functions of  $x_1, x_2, x_3$  respectively, are really fulfilled, then it must appear, that the functional values  $z_a(k_a), \zeta_{b;a}(k_a, k_b), Z_c(k_a, k_b, k_c)$ , corresponding to  $x_a(k_a), x_b(k_b), x_c(k_c)$ , satisfy the two conditions **Ia** and **Ib**. Moreover it must appear that the three magnitudes  $\gamma_{ab}$ , computed from **Ebis**, are constant, that is to say: that

$$\gamma_{ab} = \frac{A_{ab}}{Q_{ab}^2} = \text{constant (a, b = 1, 2, 3)} \quad \dots \quad \text{II}$$

Inversely, when the conditions

$$\left. \begin{aligned} z_a^2 + \zeta_{b;a}^2 &= z_b^2 + \zeta_{a;b}^2 (= q_{ab}^2), & \dots & \text{Ia} \\ Z_c^2 &= H\{z, \zeta\} - q_{ab}^2, & (a, b, c = 1, 2, 3) & \dots \text{Ib} \\ \gamma_{ab} &= \frac{A_{ab}}{Q_{ab}^2} = \text{constant}, & \dots & \text{II} \end{aligned} \right\}$$

are fulfilled, we may conclude, that *linear correlation* exists between the unimodular variables  $t_1, t_2, t_3$ , determined by

$$t_a = \frac{z_a}{\sin \psi_a} = \left( \frac{\Gamma_{aa}}{\Gamma} \right)^{1/2} \cdot z_a \quad (a = 1, 2, 3) \quad \dots \quad \text{J}$$

$t_a$  being (in virtue of  $\psi_a = \text{const.}$ ) a pure function of  $x_a$ .

If  $t_a$  is a pure function of  $x_a$  ( $a = 1, 2, 3$ ), each arrangement  $z_a, \zeta_{b;a}, Z_c$  associates with the set  $x_1, x_2, x_3$  the same point  $P(t_1, t_2, t_3)$  in the same (skew) system of coordinates  $t_1, t_2, t_3$ . Hence we obtain on the unity-sphere for each set  $x_1, x_2, x_3$  the same spherical triangle ( $\Phi$ ) of reference, thus also the same polar triangle ( $\Omega$ ). In this case the (total) coefficients of correlation of  $t_1\{x_1\}, t_2\{x_2\}, t_3\{x_3\}$  are the cosines of the (constant) sides of the polar triangle:

$$\gamma_{23} = \cos \omega_1 \quad , \quad \gamma_{31} = \cos \omega_2 \quad , \quad \gamma_{12} = \cos \omega_3.$$

If, in particular,  $t_a$  is a linear function of  $x_a, t_a = h_a(x_a - \bar{x}_a)$ , the

function  $z_a$ , too, will be a linear function of  $x_a$ , the function  $\zeta_{b;a}$  will be linear in  $x_a$  and  $x_b$ ,  $Z_c$  linear in  $x_1, x_2, x_3$ . Then the original variables  $x_1, x_2, x_3$  themselves are linearly correlated.

This statement furnishes us a rather reliable test of the linearity of the correlation between  $x_1, x_2, x_3$ .<sup>1)</sup>

Since no empirical frequency distribution is *rigorously* conform to any ideal theoretical probability scheme, and since we have to reckon with several sources of inaccuracy, as well in calculating mean values (such as  $\zeta_{2;1}(k_1 - \frac{1}{2}, k_2)$ ,  $Z_3(k_1 - \frac{1}{2}, k_2 - \frac{1}{2}, k_3)$ ) as in interpolating (e.g. to compute  $\zeta_{2;1}(k_1, k_2)$ ,  $Z_3(k_1, k_2, k_3)$ ), we cannot reasonably expect, that even in the case  $t_a \{x_a\}$  the conditions **Ia, b** and **II** will be satisfied *exactly*. We shall have to content ourselves with an *approximate* validity of these equations. Only when the quotients  $q_{a;b} : q_{b;a}$ ,  $Z_c^2 : (H \{z, \zeta\} - Q_{ab}^2)$  and the variables  $\gamma_{ab}$  show some obvious functional dependency on  $x_1, x_2, x_3$ , should we be obliged to drop the supposition  $t_a \{x_a\}$  ( $a = 1, 2, 3$ ).

If the condition **II** is not fulfilled, that is to say: if the values of  $\gamma_{bc} = \cos \omega_a$  computed for the different sets  $x_1(k_1), x_2(k_2), x_3(k_3)$ , are *unequal*, so that  $\omega_1, \omega_2, \omega_3$  are *variable magnitudes*, then we may *introduce* a *constant* polar triangle ( $\Omega$ ) with sides  $\overline{\omega}_1, \overline{\omega}_2, \overline{\omega}_3$ . In order to keep in touch, as much as possible, with the values  $\omega_1, \omega_2, \omega_3$  really found: it is preferable to choose for  $\omega_a$  an *average* value of  $\omega_a$ . We may obtain such average values by taking the averages of

$$M_1 = \text{tg } \omega_1 = \frac{B_{23}}{A_{23}} \quad , \quad M_2 = \text{tg } \omega_2 = \frac{B_{31}}{A_{31}} \quad , \quad M_3 = \text{tg } \omega_3 = \frac{B_{12}}{A_{12}}$$

according to the precept we have formerly given in S. C. I. *b* (Dutch text p. 976, English text p. 929).

$A_{23}$  and  $B_{23}$  are merely functions of  $x_2$  and  $x_3$ . So, in averaging  $M_1$ , the frequency distribution is to be treated as a two-dimensional one, its frequencies being  $Y'(k_2, k_3) = \sum_{i_1=1}^{y_1} Y(i_1, k_2, k_3)$ . Hence the "weights" introduced in averaging  $M_1 = \text{tg } \omega_1$  are built up of these two-dimensional frequencies  $Y'(k_2, k_3)$ . Likewise for averaging  $M_2$  and  $M_3$ .

From the averages  $\overline{\omega}_a$  ( $a = 1, 2, 3$ ) we compute the corresponding arcs  $\overline{\varphi}_a$  and  $\overline{\psi}_a$  ( $a = 1, 2, 3$ ).

Then we put, guided by **9bis**,

$$\left. \begin{aligned} z_1 &= \sin \overline{\psi}_1 \cdot T_1, \\ \zeta_{2;1} &= -\sin \overline{\varphi}_2 \cos \overline{\omega}_3 \cdot T_1 + \sin \overline{\varphi}_1 \cdot T_2, \\ Z_3 &= \cos \overline{\varphi}_2 \cdot T_1 + \cos \overline{\varphi}_1 \cdot T_2 + T_3. \end{aligned} \right\} \cdot \cdot \cdot \quad \mathbf{32} \left\{ \begin{array}{l} (1) \\ (2; 1) \\ (3; 2 1) \end{array} \right.$$

<sup>1)</sup> See the footnote further on, preceding equat. J in the treatment of  $n$  variables.

By confronting these equations with the equations 9bis we obtain :

$$\begin{aligned} z_1 &= \sin \bar{\psi}_1 \cdot T_1 = \sin \psi_1 \cdot t_1, \\ \zeta_{2;1} &= -\sin \bar{\varphi}_2 \cos \bar{\omega}_3 \cdot T_1 + \sin \bar{\varphi}_1 \cdot T_2 = -\sin \varphi_2 \cos \omega_3 \cdot t_1 + \sin \varphi_1 \cdot t_2, \\ Z_3 &= \cos \bar{\varphi}_2 \cdot T_1 + \cos \bar{\varphi}_1 \cdot T_2 + T_3 = \cos \varphi_2 \cdot t_1 + \cos \varphi_1 \cdot t_2 + t_3, \end{aligned}$$

whence

$$T_1 = \frac{\sin \psi_1}{\sin \varphi_1} t_1 = \frac{z_1}{\sin \psi_1} \dots \dots \dots \quad \mathbf{33}$$

$$\begin{aligned} T_2 &= \frac{1}{\sin \varphi_1} (-\sin \varphi_2 \cos \omega_3 \cdot t_1 + \sin \varphi_1 \cdot t_2 + \sin \bar{\varphi}_2 \cos \bar{\omega}_3 \cdot T_1) \\ &= \frac{1}{\sin \varphi_1} \left( -\sin \varphi_2 \cos \omega_3 + \frac{\sin \bar{\varphi}_2 \cos \bar{\omega}_3 \sin \psi_1}{\sin \varphi_1} \right) t_1 + \frac{\sin \varphi_1}{\sin \varphi_1} \cdot t_2 \quad \mathbf{34} \end{aligned}$$

$$\begin{aligned} &= \frac{1}{\sin \varphi_1 \sin \psi_1} (-\sin \varphi_2 \cos \omega_3 \sin \bar{\varphi}_2 \sin \bar{\omega}_3 + \\ &\quad + \sin \bar{\varphi}_2 \cos \bar{\omega}_3 \sin \varphi_2 \sin \omega_3) t_1 + \frac{\sin \varphi_1}{\sin \varphi_1} \cdot t_2 \end{aligned}$$

$$= \frac{\sin \varphi_2 \sin \bar{\varphi}_2 \sin (\omega_3 - \bar{\omega}_3)}{\sin \varphi_1 \sin \psi_1} t_1 + \frac{\sin \varphi_1}{\sin \varphi_1} t_2$$

$$= \frac{\sin \varphi_2 \sin (\omega_3 - \bar{\omega}_3)}{\sin \varphi_1 \sin \omega_3} t_1 + \frac{\sin \varphi_1}{\sin \varphi_1} \cdot t_2$$

$$= \frac{\sin \varphi_2 \sin (\omega_3 - \bar{\omega}_3)}{\sin \psi_2} t_1 + \frac{\sin \varphi_1}{\sin \varphi_1} t_2$$

$$= \frac{\sin \varphi_2 \sin (\omega_3 - \bar{\omega}_3)}{\sin \psi_2} \cdot \frac{1}{\sin \varphi_2 \sin \omega_3} z_1 + \frac{\sin \varphi_1}{\sin \varphi_1} \cdot \frac{1}{\sin \varphi_1 \sin \omega_3} z_2$$

or

$$\left. \begin{aligned} T_2 &= \frac{\sin (\omega_3 - \bar{\omega}_3)}{\sin \psi_2 \sin \omega_3} z_1 + \frac{1}{\sin \varphi_1 \sin \omega_3} z_2, \\ &= \frac{\sin (\omega_3 - \bar{\omega}_3)}{\sin \varphi_1 \sin \omega_3 \sin \omega_3} z_1 + \frac{1}{\sin \varphi_1 \sin \omega_3} z_2, \end{aligned} \right\} \dots \mathbf{34bis}$$

and

$$T_3 = \cos \varphi_2 \cdot t_1 + \cos \varphi_1 \cdot t_2 + t_3 - \cos \bar{\varphi}_2 \cdot T_1 - \cos \bar{\varphi}_1 \cdot T_2$$

or, by 34,

$$\begin{aligned} T_3 &= \cos \varphi_2 \cdot t_1 + \cos \varphi_1 \cdot t_2 + t_3 - \cos \varphi_2 \frac{\sin \psi_1}{\sin \varphi_1} t_1 - \\ &\quad - \cos \bar{\varphi}_1 \left( -\frac{\sin \varphi_2 \cos \omega_3}{\sin \varphi_1} + \frac{\sin \bar{\varphi}_2 \cos \bar{\omega}_3 \sin \psi_1}{\sin \varphi_1 \sin \psi_1} \right) t_1 - \cos \bar{\varphi}_1 \frac{\sin \varphi_1}{\sin \varphi_1} t_2 \end{aligned}$$

or

$$T_3 = \left\{ \frac{\cos \varphi_2 \sin \bar{\varphi}_1 + \cos \bar{\varphi}_1 \sin \varphi_2 \cos \omega_3}{\sin \bar{\varphi}_1} - \frac{(\cos \varphi_2 \sin \bar{\varphi}_1 + \cos \bar{\varphi}_1 \sin \varphi_2 \cos \omega_3) \sin \psi_1}{\sin \bar{\varphi}_1 \sin \psi_1} \right\} t_1 - \frac{\sin \varphi_1 \cos \bar{\varphi}_1 - \cos \varphi_1 \sin \bar{\varphi}_1}{\sin \bar{\varphi}_1} t_2 + t_3.$$

We now have in the "average" spherical triangle  $\bar{\Phi}_1 \bar{\Phi}_2 \bar{\Phi}_3$ :

$$\begin{aligned} \cos \bar{\varphi}_2 \sin \bar{\varphi}_1 + \cos \bar{\varphi}_1 \sin \bar{\varphi}_2 \cos \bar{\omega}_3 &= \cos \bar{\varphi}_2 \sin \bar{\varphi}_1 - \cos \bar{\varphi}_1 \sin \bar{\varphi}_2 \cos \bar{\Phi}_3 = \\ &= \sin \bar{\varphi}_3 \cos \bar{\Phi}_2 = -\sin \bar{\varphi}_3 \cos \bar{\omega}_2; \end{aligned}$$

hence the second term in the coefficient of  $t_1$  becomes:

$$+ \frac{\sin \bar{\varphi}_3 \cos \bar{\omega}_2}{\sin \bar{\varphi}_1 \sin \psi_1} \cdot \sin \psi_1 = \frac{\sin \bar{\varphi}_3 \cos \bar{\omega}_2}{\sin \bar{\varphi}_1 \sin \varphi_3 \sin \omega_2} \cdot \sin \psi_1 = \frac{\cot \bar{\omega}_2}{\sin \bar{\varphi}_1} \cdot \sin \psi_1.$$

Thus we obtain, by  $t_a = \frac{z_a}{\sin \psi_a}$  ( $a = 1, 2, 3$ ),

$$T_3 = \left( \frac{\cos \varphi_2 \sin \bar{\varphi}_1 + \cos \bar{\varphi}_1 \sin \varphi_2 \cos \omega_3}{\sin \bar{\varphi}_1 \sin \psi_1} + \frac{\cot \bar{\omega}_2}{\sin \bar{\varphi}_1} \right) z_1 - \frac{\sin (\varphi_1 - \bar{\varphi}_1)}{\sin \bar{\varphi}_1 \sin \psi_2} z_2 + \frac{1}{\sin \psi_3} z_3. \quad 35$$

The first term in the coefficient of  $z_1$  can be interpreted geometrically as follows:

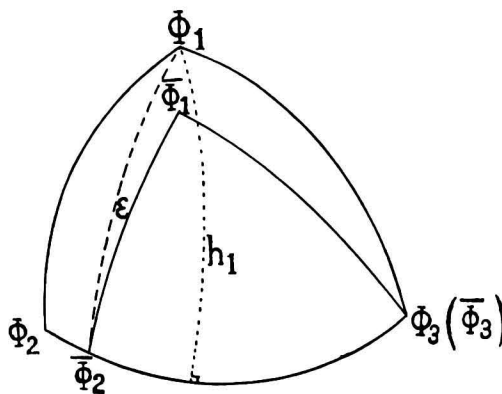


fig. 4

We locate the "average" spherical triangle so, that  $\bar{\Phi}_3$  coincides with  $\Phi_3$ , and that  $\bar{\Phi}_2 \bar{\Phi}_3$  falls along  $\Phi_2 \Phi_3$  (fig. 4).

Putting

$$\angle \overline{\Phi}_1 \overline{\Phi}_2 \overline{\Phi}_1 = \varepsilon,$$

we have in  $\triangle \overline{\Phi}_1 \overline{\Phi}_2 \overline{\Phi}_3$ :

$$\cos \varphi_2 \sin \overline{\varphi}_1 - \cos \overline{\varphi}_1 \sin \varphi_2 \cos \overline{\Phi}_3 = \sin \overline{\Phi}_2 \overline{\Phi}_1 \cdot \cos \overline{\Phi}_3 \overline{\Phi}_2 \overline{\Phi}_1,$$

or

$$\begin{aligned} \cos \varphi_2 \sin \overline{\varphi}_1 + \cos \overline{\varphi}_1 \sin \varphi_2 \cos \omega_3 &= \frac{\sin h_1}{\sin (\overline{\Phi}_2 + \varepsilon)} \cos (\overline{\Phi}_2 + \varepsilon) = \\ &= \sin \psi_1 \cot (\overline{\Phi}_2 + \varepsilon) = -\sin \psi_1 \cot (\overline{\omega}_2 - \varepsilon). \end{aligned}$$

Hence the coefficient of  $z_1$  becomes:

$$\frac{1}{\sin \varphi_1} \{-\cot (\overline{\omega}_2 - \varepsilon) + \cot \overline{\omega}_2\} = \frac{-\sin \varepsilon}{\sin \varphi_1 \sin \overline{\omega}_2 \sin (\overline{\omega}_2 - \varepsilon)} = \frac{-\sin \varepsilon}{\sin \psi_3 \sin (\overline{\omega}_2 - \varepsilon)}.$$

Thus we find for  $T_3$ :

$$T_3 = \frac{-\sin \varepsilon}{\sin \psi_3 \sin (\overline{\omega}_2 - \varepsilon)} z_1 - \frac{\sin (\varphi_1 - \overline{\varphi}_1)}{\sin \varphi_1 \sin \psi_2} z_2 + \frac{1}{\sin \psi_3} z_3. \quad \text{35bis}$$

So we have together:

$$T_1 = \frac{1}{\sin \psi_1} z_1,$$

$$T_2 = \frac{\sin (\omega_3 - \overline{\omega}_3)}{\sin \psi_2 \sin \omega_3} z_1 + \frac{1}{\sin \varphi_1 \sin \omega_3} z_2$$

$$T_3 = \frac{-\sin \varepsilon}{\sin \psi_3 \sin (\overline{\omega}_2 - \varepsilon)} z_1 - \frac{\sin (\varphi_1 - \overline{\varphi}_1)}{\sin \varphi_1 \sin \psi_2} z_2 + \frac{1}{\sin \psi_3} z_3$$

where  $\varepsilon$  is determined by

$$\cot (\overline{\omega}_2 - \varepsilon) = -\frac{\sin \overline{\varphi}_1 \cos \varphi_2 + \cos \overline{\varphi}_1 \sin \varphi_2 \cos \omega_3}{\sin \psi_1}.$$

In this way we have constructed three variables  $T_1, T_2, T_3$ , which are *linearly* correlated, having the (total) coefficients of correlation:

$$\overline{\gamma}_{23} = \cos \overline{\omega}_1, \quad \overline{\gamma}_{31} = \cos \overline{\omega}_2, \quad \overline{\gamma}_{12} = \cos \overline{\omega}_3.$$

$T_1$  is a function of only  $z_1$ , thus a pure function of  $x_1$ .

Since  $\omega_3 = \arctg \frac{B_{12}}{A_{12}}$  is built up out of  $z_1, z_2, \zeta_{2,1}, \zeta_{1,2}$ ,  $z_1$  depending only on  $x_1$ ,  $z_2$  only on  $x_2$ ,  $\zeta_{2,1}$  and  $\zeta_{1,2}$  only on  $x_1$  and  $x_2$ , we recognize  $\omega_3$  to be a function only of  $x_1$  and  $x_2$ . Thus:

$$T_2 \text{ is a function only of } x_1 \text{ and } x_2 \left( \frac{\partial T_2}{\partial x_3} = 0 \right).$$

$T_3$  however is a function of all three variables  $x_1, x_2, x_3$ .

Hence:

If the conditions Ia and Ib are fulfilled, but II is not, we may, by introducing the constant (average) magnitudes

$\overline{\omega}_1, \overline{\omega}_2, \overline{\omega}_3$ , establish a linear correlation between  $T_1\{x_1\}$ ,  $T_2\{x_1, x_2\}$ ,  $T_3\{x_1, x_2, x_3\}$ . Then the (total) coefficients of correlation are  $\gamma_{23} = \cos \omega_1$ ,  $\gamma_{31} = \cos \omega_2$ ,  $\gamma_{12} = \cos \omega_3$ .

If, the magnitudes  $\omega_1$  and  $\omega_2$  being variable, the magnitude  $\omega_3 = \arctg \frac{B_{12}}{A_{12}}$  is found to be constant, then we naturally choose this constant value for  $\overline{\omega}_3$  (fig. 5), whence  $\omega_3 - \overline{\omega}_3 = 0$ .

In this case  $T_2$  becomes a pure function of  $x_2$ .

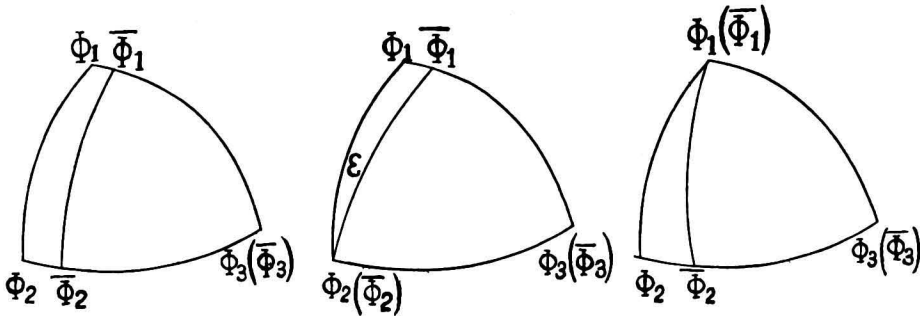


fig.5

fig.6

fig.7

If it is found that not only  $\omega_3$ , but also  $\varphi_1$  is constant (fig. 6), the expression

$$\cos \varphi_1 = \frac{\cos \omega_2 \cos \omega_3 - \cos \omega_1}{\sin \omega_2 \sin \omega_3} = \frac{A_{31} A_{12} Q_{23}^2 - A_{23} Q_{31}^2 Q_{12}^2}{B_{31} B_{12} Q_{23}^2}$$

turning out to take the same value for each set  $x_1(k_1), x_2(k_2), x_3(k_3)$ , then we have also

$$\sin \psi_2 = \sin \omega_3 \sin \varphi_1 = \text{constant.}$$

We now obtain (see fig. 6):

$$\epsilon = \overline{\Phi}_2 - \overline{\Phi}_2 = -(\omega_2 - \overline{\omega}_2) \quad , \quad \overline{\omega}_2 - \epsilon = \omega_2.$$

In this case the expression 35bis for  $T_3$  gets rid of its term with  $z_2$ , and is reduced to

$$T_3 = \frac{\sin(\omega_2 - \overline{\omega}_2)}{\sin \psi_3 \sin \omega_2} z_1 + \frac{1}{\sin \psi_3} z_3.$$

Here  $\omega_2$  is a function only of  $z_1$  and  $z_3$ ; so  $\sin \psi_3 (= \sin \varphi_1 \sin \omega_2)$  is. Therefore:  $\omega_3 = \text{const.}$  and  $\varphi_1 = \text{const.}$  furnishes  $T_1\{x_1\}$ ,  $T_2\{x_2\}$ ,  $T_3\{x_1, x_3\}$ .

If, besides  $\omega_3$ , also  $\varphi_2$  turns out to be constant, the expression

$$\cos \varphi_2 = \frac{A_{12} A_{23} Q_{31}^2 - A_{31} Q_{12}^2 Q_{23}^2}{B_{12} B_{23} Q_{31}^2}$$

taking the same value for each set  $x_1(k_1), x_2(k_2), x_3(k_3)$ , then  $\varepsilon = 0$  (see fig. 7), so that  $T_3$  drops its term with  $z_1$ , and is reduced to

$$T_3 = \frac{-\sin(\varphi_1 - \overline{\varphi_1})}{\sin \varphi_1 \sin \psi_2} z_2 + \frac{1}{\sin \psi_3} z_3.$$

Here  $\sin \psi_3 = \sin \varphi_2 \sin \omega_1 \{x_2, x_3\}$  is a function only of  $x_2$  and  $x_3$ ; however the coefficient of  $z_2$  still depends on all the variables  $x_1, x_2, x_3$ . Hence in this case  $T_3$  is still a function of all three variables  $x_a$ .

By interchanging the subscripts we may — in the same manner — arrive at analogous relations of dependency.

We may observe, that we have computed the elements of the spherical triangle of reference exclusively by means of **E**, hence out of the magnitudes  $z_a, \zeta_{b;a}$ . So we have only used these magnitudes, which are much more accurate than the magnitudes  $Z_c, 1^\circ$  because, for calculating  $Z_c$ , we have much more to do with mean values (e.g.  $z_a(k_a - \frac{1}{2})$ ) and with interpolated values (e.g.  $Z_c(k_1, k_2, k_3)$ ),  $2^\circ$  because the magnitudes  $Z_c$  (see **B(3;21)**) are computed from — generally small — one-dimensional frequency sums.

If **Ia** is satisfied, but **Ib** is not, then the 6 triplets  $z_a, \zeta_{b;a}, Z_c$  still join to each set  $x_1(k_1), x_2(k_2), x_3(k_3)$  one and the same image point  $\Pi(k_1, k_2, k_3)$  on the unity-sphere (see fig. 3), the 3 expressions  $r_{ab}^2 = q_{ab}^2 + Z_c^2$  giving however values for  $r^2$  different from the value of  $H\{z, \zeta\}$  as it is deduced solely from the magnitudes  $z_a, \zeta_{b;a}$ . We may now introduce new magnitudes  $Z'_1, Z'_2, Z'_3$ , which do satisfy **Ib**, in other words: we put:

$$Z_c'^2 = H\{z, \zeta\} - q_{ab}^2 \quad (c = 1, 2, 3) \quad . . . . \quad 36$$

and consider these  $Z'_c$  as “adjusted” values of  $Z_c$ . So we replace the empirical triplets  $z_a, \zeta_{b;a}, Z_c$ , computed from the *given* frequency distribution, by the triplets  $z_a, \zeta_{b;a}, Z'_c$ . By computing back the frequency distribution corresponding to the *adjusted* triplets, this frequency distribution will appear to be somewhat different from the given one. Then we must judge whether the discrepancy thus found between the given and the computed frequency distribution may be considered as a small accidental deviation.

We are inclined to adjust in the first place the magnitudes  $Z_c$ , these being subject to the greatest uncertainty. It is, on the other hand, the great uncertainty of  $Z_c$  which makes the frequency distribution rather insensible to an alteration of  $Z_c$ .

When we, nevertheless, in computing back the frequency distribution, desire to get back *exactly* the *given* frequencies, we may succeed by keeping one of the magnitudes  $Z_c$ , for instance  $Z_3$ , unaltered. <sup>1)</sup> Then

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<sup>1)</sup> We shall preferably keep that magnitude  $Z_c$ , which satisfies **Ib** best



the triplets  $z_1, \zeta_{2;1}, Z_3$  and  $z_2, \zeta_{1;2}, Z_3$  will give back exactly the original frequency distribution.

In this case it is necessary to take for  $r^2$ :

$$r'^2 = r_{12}^2 = q_{12}^2 + Z_3^2 \dots \dots \dots 37$$

Hence  $H \{z, \zeta\}$  must be adjusted to this value  $r_{12}^2$ . Thus some of the magnitudes  $z_3, \zeta_{1;3}, \zeta_{3;1}, \zeta_{2;3}, \zeta_{3;2}$  must certainly undergo alteration. We naturally prefer to leave the comparatively accurate magnitude  $z_3$  unaltered. As the two conditions

$$\zeta_{2;3}^2 - \zeta_{3;2}^2 = z_2^2 - z_3^2, \quad \zeta_{3;1}^2 - \zeta_{1;3}^2 = z_3^2 - z_1^2,$$

deriving from Ia, must in any case be satisfied, we must vary simultaneously either  $\zeta_{2;3}$ , and  $\zeta_{3;2}$ , or  $\zeta_{3;1}$  and  $\zeta_{1;3}$ .

Leaving, for instance,  $\zeta_{3;1}$  and  $\zeta_{1;3}$  unaltered, the alteration required in adjusting  $H \{z, \zeta\}$  will be effectuated solely on  $\zeta_{2;3}$  and  $\zeta_{3;2}$ . The magnitudes  $\zeta_{3;1}$  and  $\zeta_{1;3}$  keeping their values and  $q_{31}$  also doing so, we must first adjust  $Z_2$  into  $Z'_2$ , determined by

$$Z_2'^2 = r'^2 - q_{31}^2 \dots \dots \dots 38$$

When we, in applying the equations D and E, interchange the spherical triangle ( $\Phi$ ) and its polar triangle ( $\Omega$ ), and take into account, that, in interchanging ( $\Phi$ ) and ( $\Omega$ ), the magnitude  $z_a$  (projection of  $OP$  on  $O\Omega_a$ ) passes into the magnitude  $Z_a$  (projection of  $OP$  on  $O\Phi_a$ ), the magnitude  $\zeta_{b;a}$  (projection of  $OP$  on the line  $O\Psi_{b;a}$  of intersection of  $O\varphi_a$  and  $O\omega_c$ ) passing into  $\zeta_{b;c}$  (projection of  $OP$  on the line  $O\Psi_{b;c}$  of intersection of  $O\omega_a$  and  $O\varphi_c$ ) (see fig. 1 and 3), then we obtain from

$$\text{tg } \omega_1 = \frac{z_2 \zeta_{2;3} + z_3 \zeta_{3;2}}{z_2 z_3 - \zeta_{3;2} \zeta_{2;3}};$$

$$\text{tg } \varphi_1 = \frac{Z_2 \zeta_{2;1} + Z_3 \zeta_{3;1}}{Z_2 Z_3 - \zeta_{3;1} \zeta_{2;1}} \dots \dots \dots 39$$

Replacing in this expression  $Z_2$  by  $Z'_2$ , we obtain an adjusted value  $\varphi'_1$  of  $\varphi_1$ , determined by

$$\text{tg } \varphi'_1 = \frac{Z'_2 \zeta_{2;1} + Z_3 \zeta_{3;1}}{Z'_2 Z_3 - \zeta_{3;1} \zeta_{2;1}} \dots \dots \dots 39\text{bis}$$

The arcs  $\omega_2$  and  $\omega_3$  remaining unaltered, the arc  $\omega_1$ , which depends on  $\zeta_{2;3}$  and  $\zeta_{3;2}$ , undergoes alteration. Here the formula:

$\cos \omega_1 = \cos \omega_2 \cos \omega_3 - \sin \omega_2 \sin \omega_3 \cos \varphi_1$  furnishes, if applied to the adjusted values:

$$\cos \omega'_1 = \cos \omega_2 \cos \omega_3 - \sin \omega_2 \sin \omega_3 \cos \varphi'_1 \dots \dots \dots 40$$

The relation

$$q_{23}^2 = \sin^2 \varphi_3 \cdot t_2^2 - 2 \cos \omega_1 \sin \varphi_2 \sin \varphi_3 \cdot t_2 t_3 + \sin^2 \varphi_2 \cdot t_3^2 = \frac{z_2^2 - 2 \cos \omega_1 \cdot z_2 z_3 + z_3^2}{\sin^2 \omega_1}$$

enables us to compute the new value  $q'_{23}$  of  $q_{23}$ , viz.:

$$q'_{23} = \frac{z_2^2 - 2 \cos \omega'_1 \cdot z_2 z_3 + z_3^2}{\sin^2 \omega'_1} \dots \dots \dots \mathbf{41}$$

Putting:

$$q'_{23} = q_{23}^2 + \delta, \dots \dots \dots \mathbf{42}$$

the equations

$$\zeta'_{2;3} = \zeta_{2;3}^2 + \delta, \quad \zeta'_{3;2} = \zeta_{3;2}^2 + \delta \dots \dots \dots \mathbf{43}$$

furnish the adjusted values  $\zeta'_{2;3}$  and  $\zeta'_{3;2}$  of  $\zeta_{2;3}$  and  $\zeta_{3;2}$ .

We naturally choose the average value  $\omega'_1$  of the new values  $\omega'_1$  for the constant  $\omega_1$  in the formulae **K** for  $T_1, T_2, T_3$ . Here we use the *adjusted* spherical triangles  $(\Phi')$  and  $(\Omega')$ , determined by the unaltered sides  $\omega_2$  and  $\omega_3$ , and by the adjusted side  $\omega'_1$ . This latter must be computed from **37, 38, 39bis, 40**, thus according to the scheme:

$$\left. \begin{aligned} r'^2 &= r_{12}^2 = q_{12}^2 + Z_3^2, \\ Z_2'^2 &= r'^2 - q_{31}^2, \\ \text{tg } \varphi'_1 &= \frac{Z_2' \zeta_{2;1} + Z_3 \zeta_{3;1}}{Z_2' Z_3 - \zeta_{3;1} \zeta_{2;1}} \\ \cos \omega'_1 &= \cos \omega_2 \cos \omega_3 - \sin \omega_2 \sin \omega_3 \cos \varphi'_1 \end{aligned} \right\} \dots \dots \dots \mathbf{L}$$

If none of the conditions **Ia, b, II** is satisfied, then some of the magnitudes  $z_a, \zeta_{b;a}$  must also be adjusted. Unless we desire, intentionally, to keep at least one of the magnitudes  $Z_c$ , we prefer to determine  $\omega_1, \omega_2, \omega_3$  again by the equations **Ebis**, in which  $A_{ab}$  and  $B_{ab}$  are constructed out of the *unadjusted* values  $z_a, z_b, \zeta_{b;a}, \zeta_{a;b}$ .

Now we must make  $q_{b;a}$  agree with  $q_{a;b}$ .

Referring to S. C. II, p. 408, we observe that the most preferable methods of adjustment are the following two:

a.: we keep  $z_a$  and  $z_b$ , derive  $\omega_c$  from **Ebis**, and adjust the values  $\zeta_{b;a}$  and  $\zeta_{a;b}$  to the values  $\zeta'_{b;a}$  and  $\zeta'_{a;b}$ , which are, in virtue of **9ter (b; a)** determined by

$$\zeta'_{b;a} = \frac{z_b - \cos \omega_c \cdot z_a}{\sin \omega_c}, \quad \zeta'_{a;b} = \frac{z_a - \cos \omega_c \cdot z_b}{\sin \omega_c} \dots \dots \dots \mathbf{44}$$

(see fig. 8).

So we adjust  $q_{b;a}$  and  $q_{a;b}$  to  $q'_{b;a} = q'_{a;b} = q'_{ab}$ , determined by

$$q'_{ab}{}^2 = z_a^2 + \zeta'_{b;a}{}^2 = z_b^2 + \zeta'_{a;b}{}^2 = \frac{z_a^2 - 2 \cos \omega_c \cdot z_a z_b + z_b^2}{\sin^2 \omega_c} \dots \dots \dots \mathbf{45}$$

Then the function  $r^2 = H\{z, \zeta\}$ , being (by **31**) built up of the magni-

tudes  $z_a$  and  $\omega_a$  ( $a = 1, 2, 3$ ) may be computed immediately by means of the values  $\omega_1, \omega_2, \omega_3$  calculated from **Ebis**.

So we have evidently:

$$H\{z, \zeta'\} = H\{z, \zeta\}.$$

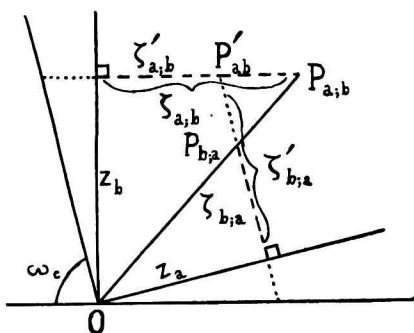


fig.8

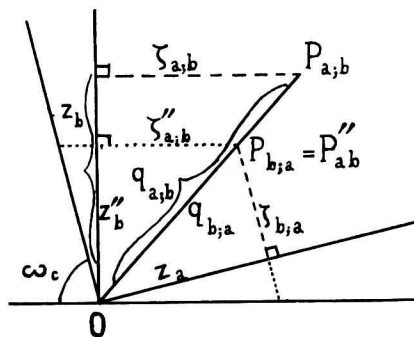


fig.9

Then the magnitudes  $Z_c$  are adjusted into  $Z'_c$  according to 36.

The spherical triangles  $(\Phi)$  and  $(\Omega)$  remaining unaltered, we derive the unimodular, linearly correlated variables  $T_1, T_2, T_3$  still from the same equations **K**, we formerly constructed.

In computing back the frequency distribution out of one of the triplets  $z_a, \zeta'_{b;a}, Z'_c$ , we must be prepared to find deviations from the original frequency distribution. The advantage of this method is, however, that we make as much use as possible of the magnitudes  $z_1, z_2, z_3$ , which are the most accurate ones.

b. We keep  $z_a$  and  $\zeta_{b;a}$ , thus also  $q_{b;a}^2 = z_a^2 + \zeta_{b;a}^2$ , and derive  $\omega_c$  still from **Ebis** (fig. 9). The magnitude  $z_b$  however must now be replaced by  $z''_b$ , which is, in virtue of **9ter** (**b; a**), determined by

$$z''_b = \cos \omega_c \cdot z_a + \sin \omega_c \cdot \zeta_{b;a} \quad \dots \quad \mathbf{M}$$

By means of **E** we may transform this expression into:

$$z''_b = \frac{z_a A_{ab} + \zeta_{b;a} B_{ab}}{q_{b;a} \cdot q_{a;b}} = \frac{z_a (z_a z_b - \zeta_{b;a} \zeta_{a;b}) + \zeta_{b;a} (z_a \zeta_{a;b} + z_b \zeta_{b;a})}{q_{b;a} \cdot q_{a;b}} = \frac{z_b \cdot q_{b;a}^2}{q_{b;a} \cdot q_{a;b}}$$

or

$$z''_b = \frac{q_{b;a}}{q_{a;b}} \cdot z_b \quad \dots \quad \mathbf{Mbis}$$

Therefore, if we wish to keep, for instance,  $z_1$  and  $\zeta_{2;1}$ , we must replace  $z_2$  by  $z''_2 = \frac{q_{2;1}}{q_{1;2}} \cdot z_2$ , and substitute this value  $z''_2$  for  $z_2$  in the equations **K**. The method (b) is to be followed, if we desire, intentionally, to get back the original frequency distribution out of one of the triplets. In this case we have, of course, to leave also  $Z_c$  (i.c.  $Z_3$ )

unaltered, and to apply the corresponding method of adjustment, however on the understanding, that we must take for  $r^2$  the value:

$$r''^2 = r_{2;1}^2 = q_{2;1}^2 + Z_3^2 = z_1^2 + \zeta_{2;1}^2 + Z_3^2, \dots \dots \dots 46$$

so that also the values of  $Z'_2, \varphi'_1, \omega'_1$  turn out to be different from those derived from L.

If we desire to keep  $\zeta_{3;1}$  and  $\zeta_{1;3}$ , we must, while retaining  $\omega_2$ , replace these magnitudes, according to 44, by

$$\zeta''_{3;1} = \frac{z_3 - \cos \omega_2 \cdot z_1}{\sin \omega_2}, \quad \zeta''_{1;3} = \frac{z_1 - \cos \omega_2 \cdot z_3}{\sin \omega_2} \dots \dots \dots 47$$

Then we must take for  $q_{31}^2$  (see 41):

$$q_{31}''^2 = \frac{z_3^2 - 2 \cos \omega_2 \cdot z_3 z_1 + z_1^2}{\sin^2 \omega_2} \dots \dots \dots 48$$

So we obtain for  $Z_2$  the adjusted value  $Z''_2$ , determined by

$$Z_2''^2 = r''^2 - q_{31}''^2 \dots \dots \dots 49$$

By means of the value  $\zeta''_{3;1}$ , found from 47, and the value  $Z''_2$ , derived from 49, we now compute the adjusted value  $\varphi''_1$  from

$$\text{tg } \varphi''_1 = \frac{Z_2'' \zeta_{2;1} + Z_3 \zeta''_{3;1}}{Z_2'' Z_3 - \zeta''_{3;1} \zeta_{2;1}} \dots \dots \dots 50$$

At last we find the adjusted value  $\omega''_1$  from

$$\cos \omega''_1 = \cos \omega_2 \cos \omega_3 - \sin \omega_2 \sin \omega_3 \cos \varphi''_1 \dots \dots \dots 51$$

Then we determine the average value  $\overline{\omega''_1}$  of this adjusted value  $\omega''_1$ , after which we substitute this value  $\overline{\omega''_1}$  for the constant  $\omega_1$  in the formulae K. Hence, in the actual case ( $z_1, \zeta_{2;1}, Z_3$  being kept) the formulae K are altered in such a way that we replace  $z_2$  by  $z''_2$  (determined by M (Mbis)) and  $\overline{\omega_1}$  by  $\overline{\omega''_1}$  (determined by 46–51). So the adjustment is to be effectuated according to the scheme:

$$\left. \begin{aligned} z''_2 &= \cos \omega_3 \cdot z_1 + \sin \omega_3 \cdot \zeta_{2;1} = \frac{q_{2;1}}{q_{1;2}} \cdot z_2 \\ r''^2 &= r_{2;1}^2 = z_1^2 + \zeta_{2;1}^2 + Z_3^2 \\ \zeta''_{3;1} &= \frac{z_3 - \cos \omega_2 \cdot z_1}{\sin \omega_2} \\ q_{31}''^2 &= \frac{z_3^2 - 2 \cos \omega_2 \cdot z_3 z_1 + z_1^2}{\sin^2 \omega_2} \\ Z_2''^2 &= r''^2 - q_{31}''^2 \\ \text{tg } \varphi''_1 &= \frac{Z_2'' \zeta_{2;1} + Z_3 \zeta''_{3;1}}{Z_2'' Z_3 - \zeta''_{3;1} \zeta_{2;1}} \\ \cos \omega''_1 &= \cos \omega_2 \cos \omega_3 - \sin \omega_2 \sin \omega_3 \cos \varphi''_1 \end{aligned} \right\} \dots \dots \dots \text{N}$$

Of course the other elements  $\varphi_a, \psi_a$  of the spherical triangles occurring in **K**, are to be recomputed from  $\omega''_1, \omega_2, \omega_3$ . Thereby they undergo alteration, and so do the constants  $\varphi_a, \psi_a$  derived from  $\omega''_1, \omega_2, \omega_3$ .

**Summary of the Treatment of Skew Correlation between Three Variables:**

From the given frequency distribution **A** the magnitudes  $z_a, \zeta_{b;a}, Z_c$  are calculated by means of **B**.

The equations **C (Cbis)**, **D, E (Ebis)**, **F, G, H** define the magnitudes  $q_{b;a}, Q_{ab}, A_{ab}, B_{ab}, M_{ab}, \gamma_{ab}, r^2 = H\{z, \zeta\}$ , by means of which we can formulate the conditions **Ia, b, II**:

$$\begin{aligned}
 q_{b;a} &= q_{a;b} (= q_{ab}) & (a, b = 1, 2, 3) \dots \dots \dots & \text{Ia} \\
 Z_c^2 &= H\{z, \zeta\} - q_{ab}^2 & (c = 1, 2, 3) \dots \dots \dots & \text{Ib} \\
 \frac{\sqrt{1 - \gamma_{ab}^2}}{\gamma_{ab}} &= M_{ab} = \frac{B_{ab}}{A_{ab}} = \text{constant} & \dots \dots \dots & \text{II}
 \end{aligned}$$

If both **Ia, b** and **II** are satisfied for each set  $x_1(k_1), x_2(k_2), x_3(k_3)$ , there exists *linear* correlation between the unimodular variables  $t_1\{x_1\}, t_2\{x_2\}, t_3\{x_3\}$ , determined by the equations **J**, the total coefficients of correlation being  $\gamma_{23}, \gamma_{31}, \gamma_{12} \left( \frac{\gamma_{ab}}{A_{ab}} > 0 \right)$ .

If **Ia, b** are satisfied, but **II** is not, then we can establish linear correlation between the (unimodular) variables  $T_1\{x_1\}, T_2\{x_1, x_2\}, T_3\{x_1, x_2, x_3\}$  given by the equations **K**, the total coefficients of correlation being  $\gamma_{23}, \gamma_{31}, \gamma_{12}$ .

If **Ib** is not satisfied, then we may, by abandoning — if necessary — all the magnitudes  $Z_c$ , keep the variables  $T_1\{x_1\}, T_2\{x_1, x_2\}, T_3\{x_1, x_2, x_3\}$  determined by **K**. If, however, we insist upon keeping  $Z_3$  intact, then either  $\omega_1$  or  $\omega_2$  must be altered. When we keep  $\omega_2$ , the magnitude  $\omega_1$  must be adjusted to the magnitude  $\omega'_1$ , to be computed by means of the equations **L**. Then this  $\omega'_1$  and its average  $\overline{\omega'_1}$  take the place of  $\omega_1$  and  $\overline{\omega_1}$  in the formulae **K**.

If even **Ia** is not fulfilled, we may, by abandoning — if necessary — all the magnitudes  $\zeta_{b;a}$ , but keeping the magnitudes  $\omega_1, \omega_2, \omega_3$ , computed from **Ebis**, retain, even in this case, the variables  $T_1\{x_1\}, T_2\{x_1, x_2\}, T_3\{x_1, x_2, x_3\}$  defined by **K**.

If, on the contrary, we wish to keep  $\zeta_{2;1}$  intact, we must replace  $z_2$  by the magnitude  $z''_2$ , determined by **M (Mbis)**. If we desire to keep not only  $\zeta_{2;1}$  but also  $Z_3$  unaltered, we must replace the magnitudes  $z_2$  and  $\omega_1$  by the magnitudes  $z''_2$  and  $\omega''_1$  determined by **N**, and reconstruct the formulae **K**, using those adjusted values  $z''_2, \omega''_1$  and the average  $\overline{\omega''_1}$ . Then the arcs  $\varphi_a, \psi_a$  ( $a = 1, 2, 3$ ) must first be computed from  $\omega''_1, \omega_2, \omega_3$ , and their averages  $\varphi_a, \psi_a$  from  $\omega''_1, \omega_2, \omega_3$ .

(To be continued),