Mathematics. - On the differential geometry of contact transformations. By Hwa-Chung Lee. (Communicated by Prof. J. A. Schouten).
(Communicated at the meeting of September 25, 1937.)

## 1. Introduction.

In two papers of these Proceedings Schouten ${ }^{1}$ ) has shown that a doubly homogeneous contact transformation in the $2 n+2$ variables $x^{\lambda}, p_{\lambda}(x, \lambda, \ldots, \tau=0,1, \ldots, n)$ :

$$
\begin{equation*}
x^{\chi^{\prime}}=x^{x^{\prime}}\left(x^{\star}, p_{\lambda}\right), p_{\lambda^{\prime}}=p_{\lambda^{\prime}}\left(x^{x}, p_{\lambda}\right), . \tag{1}
\end{equation*}
$$

where $x^{\nu^{\prime}}$ and $p_{\lambda^{\prime}}$ are homogeneous functions of degrees $(1,0)$ and $(0,1)$ in $x^{2}, p_{2}$ respectively, can always be modified (without changing the geometrical meaning) in such a way that these functions satisfy the equations

$$
\left.\begin{array}{ll}
p_{x^{\prime}} & \frac{\partial x^{x^{\prime}}}{\partial x^{\lambda}}=p_{\lambda},  \tag{2}\\
p_{x^{\prime}} \frac{\partial x^{x^{\prime}}}{\partial p_{x}}=0 \\
x^{\lambda^{\prime}} \frac{\partial p_{\lambda^{\prime}}}{\partial x^{\lambda}}=0, & x^{\lambda^{\prime}} \frac{\partial p_{\lambda^{\prime}}}{\partial p_{x}}=x^{x},
\end{array}\right\}
$$

for all values of $x^{\star}$ and $p_{\lambda}$, though the latter only have a geometric meaning when they are bound by the relation

$$
\begin{equation*}
p_{\lambda} x^{\lambda}=0 \tag{3}
\end{equation*}
$$

Transformations of the form (1) with the mentioned homogeneity property and satisfying (2) evidently form a group ${ }^{2}$ ) which we call $\Omega_{2 n+2}$. Let us consider the $2 n+1$ dimensional manifold with the homogeneous coordinates $x^{x}, p_{\lambda}$, endowed with the group $\Omega_{2 n+2}$ of coordinate transformations and the group $\mathfrak{F}$ of point transtormations

$$
\begin{equation*}
{ }^{\prime} x^{\star}=\varrho x^{\star}, \quad{ }^{\prime} p_{\lambda}=\varrho p_{\lambda} \tag{4}
\end{equation*}
$$

where $\varrho$ is a homogeneous function of degree $(0,0)$ in $x^{x}, p_{1}$. A point ( $x^{x} \cdot p_{\lambda}$ ) of this manifold is called an element if its coordinates satisfy the

[^0]relation (3). The totality of elements constitutes a $2 n$-dimensional space $K_{2 n}$ whose defining equation (3) is invariant under the group $\Omega_{2 n+2}$ and the group $\mathfrak{F} .^{\prime}$ It is only the elements of this space with which we shall be concerned, and the purpose of this paper is to establish a formalism for the invariant theory of the groups $\Re_{2 n+2}$ and $\mathscr{F}^{\prime}$ by constructing certain fundamental projective tensors which we shall call contact projectors.

We may get a correspondence between the space $K_{2 n}$ and another space $H_{n}$ if we regard $x^{x}, p_{2}$ as two separate sets of homogeneous coordinates, the first being homogeneous coordinates of the points of $H_{n}{ }^{3}$ ) and the second being homogeneous coordinates of the hyperplanes in the local spaces of $H_{n}$ with $\left\lfloor x^{*} 」^{4}\right.$ ) as points of contact. The elements of $K_{2 n}$ correspond to the hyperplanes in the local spaces of $H_{n}$ passing through the points of contact, owing to the relation (3). The configuration formed by a point of $H_{n}$ and a hyperplane through it is also called an element of $H_{n}$. Since the two identical elements ( $x^{x}, p_{\lambda}$ ) and ( $\varrho x^{x}, \sigma p_{\lambda}$ ) of $H_{n}$ for $\sigma \neq \varrho$ are regarded as two distinct elements of $K_{2 n}$ because of (4), the elements of $K_{2 n}$ and those of $H_{n}$ are then not in one-to-one correspondence.

## 2. Contact projectors.

Besides the index-type $x$, we introduce the bracketed index-type ( $x$ ) defined by $(x)=(n+1)+x$, so that when $x$ takes the range $0,1, \ldots, n$, $(x)$ will take the continuation range $n+1, \ldots, 2 n+1$. We shall write $x^{(\lambda)}$ for $p_{\lambda}$ so that the $2 n+2$ variables $x^{x}, p_{\lambda}$ can be written as $x^{x}, x^{(\lambda)}$ or $x^{\alpha}$, where $\alpha, \beta, \ldots, \iota$ take the two ranges of $x,(\varkappa)$ successively, that is to say, $\alpha, \beta, \ldots, \iota=0,1, \ldots, 2 n+1$.

Equations (1) and (4) may now be written

$$
\begin{align*}
\Omega_{2 n+2} & : x^{\alpha^{\prime}}=x^{\alpha^{\prime}}\left(x^{\alpha}\right),  \tag{5}\\
\mathfrak{F}^{\prime} & :  \tag{6}\\
& x^{\alpha}=\varrho x^{\alpha},
\end{align*}
$$

where $x^{\alpha^{\prime}}$ are homogeneous functions of degree one in $x^{\alpha}$ and $\varrho$ is a homogeneous function of degree zero in $\boldsymbol{x}^{\alpha}$.

The transformations (5) being supposed to be non-singular, we define contact projectors with the aid of the quantities

$$
\left.\begin{array}{l}
\mathscr{\gamma}_{\beta}^{\alpha^{\prime}}=\partial_{\beta} x^{\alpha^{\prime}}, \partial_{\beta}=\frac{\partial}{\partial x^{\beta}},  \tag{7}\\
\mathscr{H}_{\beta^{\prime}}^{\alpha}=\partial_{\beta^{\prime}} x^{\alpha}, \partial_{\beta^{\prime}}=\frac{\partial}{\partial x^{\beta^{\prime}}},
\end{array}\right\} .
$$

[^1]For example a contra- or covariant (projective) contact vector $v^{\alpha}$ or $w_{\beta}$ of degree $r$ is defined by the laws

$$
\Omega_{2 n+2}:\left\{\begin{array}{l}
v^{\alpha^{\prime}}=\mathscr{\gamma}_{\alpha}^{\alpha^{\prime}} v^{\alpha},  \tag{8}\\
w_{\beta^{\prime}}=\mathscr{\gamma}_{\beta^{\prime}}^{\beta} w_{\beta},
\end{array} \quad \mathscr{F}^{\prime}:\left\{\begin{array}{l}
v^{\alpha}=\varrho^{r} v^{\alpha}, \\
\prime \quad w_{\beta}=\varrho^{r} w_{\beta} .
\end{array}\right\}\right.
$$

Let us now write equations (2) in the form

$$
\left.\begin{array}{l}
p_{\beta}=P_{\beta^{\prime}} \mathscr{X}_{\beta}^{\beta^{\prime}},  \tag{9}\\
X_{\beta}=X_{\beta^{\prime}} \mathscr{X}_{\beta}^{\beta^{\prime}},
\end{array}\right\}
$$

where $X_{\beta}$ and $P_{\beta}$ are defined in each coordinate system by

$$
\left.\begin{array}{l}
X_{\lambda}=0, X_{(x)}=x^{x}  \tag{10}\\
P_{\lambda}=p_{\lambda}, P_{(x)}=0,
\end{array}\right\}
$$

i. e. the components of $X_{\beta}$ only involve $x^{2}$, and those of $P_{\beta}$ only involve $p_{\lambda}$. Then according to the definition (8), $X_{\beta}$ and $P_{\beta}$ are covariant contact vectors of degree one.

From (9) we have by differentiation

$$
\left.\begin{array}{l}
\partial_{[\gamma} P_{\beta]}=\partial_{\left[\gamma^{\prime}\right.} P_{\left.\beta^{\prime}\right]} \mathscr{X}_{\gamma}^{\gamma^{\prime}} \mathscr{\gamma}_{\beta}^{\beta^{\prime}},  \tag{11}\\
\partial_{[\gamma} X_{\beta]}=\partial_{\left[\gamma^{\prime}\right.} X_{\left.\beta^{\prime}\right]} \mathscr{\sim}_{\gamma}^{\gamma^{\prime}} \mathscr{\gamma}_{\beta}^{\beta^{\prime}} ;
\end{array}\right\}
$$

and from (10) we find by actual calculation

$$
\begin{equation*}
-2 \partial_{[\beta} X_{\alpha]}=\epsilon_{\beta \alpha}=+2 \partial_{[\beta} P_{\alpha]}, \tag{12}
\end{equation*}
$$

where $\epsilon_{\beta \alpha}$ is defined in each coordinate system by

$$
\begin{equation*}
\epsilon_{\lambda x}=0, \epsilon_{2(x)}=-\delta_{\lambda,}^{\prime}, \epsilon_{(\lambda) x}=\delta_{x}^{2}, \epsilon_{(\lambda)(x)}=0 . \tag{13}
\end{equation*}
$$

Hence equations (11) reduce to

$$
\begin{equation*}
\left.\epsilon_{\beta \alpha}=\epsilon_{\beta^{\prime} \alpha^{\prime}} \mathcal{X}_{\beta}^{\beta^{\prime}} \mathcal{H}_{\alpha}^{\alpha^{\prime}}{ }^{5}\right) \tag{14}
\end{equation*}
$$

from which it follows that $\epsilon_{\beta \alpha}$ is a covariant contact bivector of degree zero.

The determinant of $\epsilon_{\beta \alpha}$ being unity, we can define a contravariant contact bivector $\epsilon^{\alpha \beta}$ of degree zero by the relation

$$
\epsilon^{\alpha \gamma} \epsilon_{\beta \gamma}=\mathscr{X}_{\beta}^{\alpha}\left\{\begin{array}{l}
=1 \text { for } \alpha=\beta,  \tag{15}\\
=0 \text { for } \alpha \neq \beta .
\end{array}\right\} .
$$

[^2]From this we find that the components of $\epsilon^{\alpha \beta}$ are

$$
\begin{equation*}
\epsilon^{\times \lambda}=0, \epsilon^{\times(\lambda)}=-\delta_{\lambda,}^{x}, \epsilon^{(x) \lambda}=+\delta_{x}^{\lambda}, \epsilon^{(x)(\lambda)}=0 . \tag{16}
\end{equation*}
$$

We shall use the two $\epsilon$ 's to raise and lower indices, convening that in lowering indices we sum with respect to the left index of $\epsilon_{\beta \alpha}$, while in raising indices we sum with respect to the right index of $\epsilon^{\alpha \beta}$, and shall identify those objects which are obtained from one another by raising and lowering indices.

With this understanding we obtain from (10) two contravariant contact vectors of degree one:

$$
\begin{align*}
X^{\star} & =-x^{\star}, & X^{(\lambda)} & =0,  \tag{17}\\
P^{\star} & =0, & & P^{(\lambda)}
\end{align*}=+p_{\star},
$$

whose laws of transformation

$$
\left.\begin{array}{rl}
X^{\alpha^{\prime}} & =X^{\alpha} \mathscr{H}_{\alpha}^{\alpha^{\prime}},  \tag{18}\\
P^{\alpha^{\prime}} & =P^{\alpha} \mathscr{H}_{\alpha}^{\alpha^{\prime}} .
\end{array}\right\}
$$

express analytically the fact that $x^{x^{\prime}}$ and $p_{\lambda^{\prime}}$ are homogeneous of degrees $(1,0)$ and $(0,1)$ in $x^{x}, p_{\lambda}$ respectively. Thus the homogeneity property of the group $\Omega_{2 n+2}$ is merely a consequence of (9), that is to say, a consequence of (2).

From (17) the two contact vectors $X^{\alpha}, P^{\alpha}$ are connected by the relation

$$
\begin{equation*}
x^{\alpha}=P^{\alpha}-X^{\alpha} \tag{19}
\end{equation*}
$$

and hence the coordinates $x^{\alpha}$ themselves are the components of a contravariant contact vector of degree one whose law of transformation simply expresses that the $x^{\alpha^{\prime}}$ are homogeneous of degree one in the $x^{\alpha}$.

From (10) the contact vector defined by

$$
\begin{equation*}
q_{\beta}=p_{\beta}+X_{\beta}, \quad\left(q_{\lambda}=p_{\lambda}, q_{(x)}=x^{\chi}\right) . \tag{20}
\end{equation*}
$$

is seen to be the gradient of the contact scalar $q=p_{\mu} x^{\mu}$ :

$$
\begin{equation*}
q_{\beta}=\partial_{\beta} q \tag{21}
\end{equation*}
$$

Corresponding to the relations (19), (20) we have by raising and lowering suffices,

$$
\left.\begin{array}{l}
x_{\beta}=P_{\beta}-X_{\beta},  \tag{22}\\
q^{\alpha}=P^{\alpha}+X^{\alpha} .
\end{array}\right\}
$$

It may be remarked that by taking the determinant on both sides of (14) we have $\operatorname{Det}\left(\mathcal{X}_{\alpha}^{\alpha^{\prime}}\right)= \pm 1$ and hence if we confine ourselves to the subgroup of $\Omega_{2 n+2}$ for which $\operatorname{Det}\left(\mathcal{X}_{\alpha}^{\alpha^{\prime}}\right)=+1$, the difference between contact projectors and projector-densities vanishes.

## 3. The contact connection.

Let a symmetric projective connection $\Pi_{\gamma \beta}^{\alpha}$ of degree -1 be introduced in the $K_{2 n}$ in such a way that we have

$$
\begin{equation*}
\nabla_{\gamma} \epsilon_{\beta \alpha}=\partial_{\gamma} \epsilon_{\beta \alpha}-\Pi_{\gamma \beta}^{\delta} \epsilon_{\partial \alpha}-\Pi_{\gamma \alpha}^{\delta} \epsilon_{\beta \delta}=0 \tag{23}
\end{equation*}
$$

If we multiple this by $\epsilon^{\beta \alpha}$, we obtain, on account of (15) and of the constancy of $\epsilon_{\beta \alpha}$,

$$
\begin{equation*}
\Pi_{\gamma \alpha}^{\alpha}=0 \tag{24}
\end{equation*}
$$

Hence there is no difference between the covariant differentiation of a contact projector and that of a projector-density.

From (15) and (23) we also have

$$
\begin{equation*}
\nabla_{\gamma} \epsilon^{\alpha \beta}=0 \tag{25}
\end{equation*}
$$

The contact connection $\Pi_{\gamma \beta}^{\alpha}$ may be further particularized in the following way. The quantities $x^{\alpha}$ and $q_{\beta}$ may be regarded as fundamental in the sense that in terms of them the two other quantities $X^{\alpha}, P^{\alpha}$ (and also $X_{\beta}, P_{\beta}$ ) can be expressed owing to the relations (19), (20) and (22). We have then only to consider the two contact projectors

$$
\begin{equation*}
Q_{\gamma}^{\alpha}=\nabla_{\gamma} x^{\alpha}, \quad R_{\gamma \beta}=\nabla_{\gamma} q_{\beta} . \tag{26}
\end{equation*}
$$

Since $q=p_{\lambda} x^{2}$ is a quadratic function of $x^{\alpha}$, we have by (21) the condition of homogeneity

$$
\begin{equation*}
x^{\alpha} q_{\alpha}=2 q \tag{27}
\end{equation*}
$$

from which a relation between the quantities defined by (26) may be obtained by covariant differentiation,

$$
\begin{equation*}
Q_{\gamma}^{\alpha} q_{\alpha}+R_{\gamma \alpha} x^{\alpha}=2 q_{\gamma} \tag{28}
\end{equation*}
$$

The simplest solution of this equation is

$$
\begin{equation*}
\mathrm{Q}_{\gamma}^{\alpha}=2 \not \gamma_{\gamma}^{\alpha} . \quad R_{\gamma \beta}=0, \tag{29}
\end{equation*}
$$

which means that the connection $\Pi_{\gamma \beta}^{\alpha}$ is a point displacement. ${ }^{6}$ )
In terms of the two quantities $Q_{\beta}{ }^{\alpha}, R_{\beta \alpha}$, the derivatives of the four vectors $x, q, X, P$ are

$$
\left.\begin{array}{ll}
\nabla_{\gamma} x^{\alpha}=Q_{\gamma}^{\alpha}, & \nabla_{\gamma} x_{\beta}=Q_{\gamma \beta},  \tag{30}\\
\nabla_{\gamma} q^{\alpha}=R_{\gamma}^{\alpha}, & \nabla_{\gamma} q_{\beta}=R_{\gamma \beta}, \\
\nabla_{\gamma} X^{\alpha}=\frac{1}{2}\left(R_{\gamma}^{\cdot \alpha}-Q_{\gamma}^{\cdot \alpha}\right), & \nabla_{\gamma} X_{\beta}=\frac{1}{2}\left(R_{\gamma \beta}-Q_{\gamma \beta}\right), \\
\nabla_{\gamma} P^{\alpha}=\frac{1}{2}\left(R_{\gamma}^{\cdot \alpha}+Q_{\gamma}^{\cdot \alpha}\right), & \nabla_{\gamma} P_{\beta}=\frac{1}{2}\left(R_{\gamma \beta}+Q_{\gamma \beta}\right),
\end{array}\right\}
$$

[^3]which reduce in the particular case (29) to
\[

\left.$$
\begin{array}{ll}
\nabla_{\gamma} x^{\alpha}=2 \not \gamma_{\gamma}^{\alpha}, & \nabla_{\gamma} x_{\beta}=2 \epsilon_{\gamma \beta},  \tag{31}\\
\nabla_{\gamma} q^{\alpha}=0, & \nabla_{\gamma} q_{\beta}=0 \\
\nabla_{\gamma} X^{\alpha}=-\not \gamma_{\gamma}^{\alpha}, & \nabla_{\gamma} X_{\beta}=-\epsilon_{\gamma \beta}, \\
\nabla_{\gamma} P^{\alpha}=+\not \gamma_{\gamma}^{\alpha}, & \nabla_{\gamma} P_{\beta}=+\epsilon_{\gamma \beta} .
\end{array}
$$\right\}
\]

In conclusion the author wishes to thank Dr. J. HaAntjes and Professor J. A. Schouten whose criticisms on this note have led to several improvements

Mathematics. - Conformal representations of an n-dimensional euclidean space with a non-definite fundamental form on itself. By J. Hanntjes. (Communicated by Prof. W. v. D. Woude).
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## Introduction.

It is wellknown that every real conformal point transformation in an $n$-dimensional $(n>2)$ euclidean space $\left(R_{n}\right)$ with a definite fundamental quadratic form can be brought about by a motion and an inversion or a dilatation ${ }^{1}$ ). This theorem, which for $n=3$ is due to Liouville and is called Liouville's theorem, does not hold in a euclidean space with a fundamental form which is not definite.

The problem with which we are here concerned is to find the extension of the above theorem to a euclidean manifold, the fundamental form of which is not definite. This leads to a new class of conformal transformations (formula (26)). If $\mathfrak{M}$ denotes this class, then, as we shall see, the extension of LIOUVILLE's theorem may be formulated as follows. Every real conformal point transformation in an $R_{n}(n>2)$ is composed of a motion and a transformation $T$, where $T$ is either a dilatation or an inversion or a transformation belonging to the class $\mathfrak{M}$.

It will appear that every transformation belonging to $\mathfrak{M}$ is the product of two inversions. Thus the following theorem holds in any euclidean space. The inversions and motions in an $R_{n}$ define together the conformal group of point transformations.

## § 1. Conformal transformations of the fundamental tensor.

Let $a_{\lambda x}$ be the fundamental tensor in an $n$-dimensional RiEmanNian space $V_{n}$. A transformation of the form

$$
\begin{equation*}
' a_{k x}=\sigma a_{i x} \tag{1}
\end{equation*}
$$

${ }^{1}$ ) S. Lie, Ueber Komplexe, insbesondere Linien- und Kugelkomplexe, mit Anwendung auf die Theorie partieller Differentialgleichungen, Math. Ann., 5 (1872) p. 184.


[^0]:    ${ }^{1}$ ) J. A. Schouten: Zur Differentialgeometrie der Gruppe der Berührungstransformationen. Vol. 40 (1937), 100-107, 236-245.
    ${ }^{2}$ ) Strictly speaking this is a pseudogroup. Cf. O. Veblen and J. H. C. Whitehead : The Foundations of Differential Geometry, Cambr. Tracts, p. 38.

[^1]:    ${ }^{3}$ ) $H_{n}$ is the generalized projective space with VAN DANTZIGs homogeneous coordinates. Cf. D. van Dantzig: Theorie des projektiven Zusammenhangs $n$-dimensionaler Räume. Math. Ann. 106 (1932), 400.
    ${ }^{4}$ ) $\left\lfloor x^{x}\right\rfloor$ means $x^{x}$ determined up to an arbitrary factor.

[^2]:    ${ }^{5}$ ) This is equivalent to the LAGRANGE parenthesis condition for contact transformations. Similarly the transformation law of $\epsilon^{\alpha \beta}$ defined by (15) is the equivalent of the Poisson parenthesis condition.

[^3]:    ${ }^{6}$ ) J. A. Schouten und J. HaAntjes: Zur allgemeinen projektiven Differentialgeometrie. Compositio Math. 3 (1936), p. 23.

